

## Article

# Stability, Electronic Structure and Thermodynamic Properties of Nanostructured MgH<sub>2</sub> Thin Films

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**Abstract:** Magnesium is an attractive hydrogen storage candidate due to its high gravimetric and volumetric storage capacities (7.6 wt.% and 110 gH<sub>2</sub>/l, respectively). Unfortunately, its use as a storage material for hydrogen is hampered by the high stability of its hydride, its high dissociation temperature of 573–673 K and its slow reaction kinetics. In order to overcome those drawbacks, an important advancement toward controlling the enthalpy and desorption temperatures of nanostructured MgH<sub>2</sub> thin films via stress/strain and size effects is presented in this paper, as the effect of the nano-structuring of the bulk added to a biaxial strain on the hydrogen storage properties has not been previously investigated. Our results show that the formation heat and decomposition temperature correlate with the thin film's thickness and strain/stress effects. The instability created by decreasing the thickness of MgH<sub>2</sub> thin films combined with the stress/strain effects induce a significant enhancement in the hydrogen storage properties of MgH<sub>2</sub>.

**Keywords:** DFT calculations; hydrogen storage; MgH<sub>2</sub> thin films; strain; stress; size

## 1. Introduction

The development of power storage methods and technologies is considered one of the most important efforts of this decade, in which hydrogen has become an important element among the clean and renewable sources of energy. Its combustion produces around three times more energy compared to fossil fuels and without any greenhouse gas emissions [1]. This has led many researchers and industrialists to consider hydrogen the key to the global energy revolution. Nevertheless, the safe, efficient and secure storage and transportation of this element remain the major obstacles facing its wide use and commercialization. Additionally, there are three technologies available for hydrogen storage, which are storing the gas under high pressure, storing it in liquid form at low temperature or storing it in solid materials. The first two methods have disadvantages relating to the difficulties in forming the two “high pressure and low temperature” conditions. On the other hand, solid materials have recently been used for hydrogen storage by way of hydrogen atoms being absorbed and stored in metallic matrixes, such as intermetallic materials, organic compounds, metal hydrides, porous materials and complex hydrides [2–9].

Due to the increase in interest in the development of hydrogen storage properties, our study focuses on the thermodynamic properties of magnesium hydride ( $\text{MgH}_2$ ) in both its bulk and thin film forms, with or without the application of strain.  $\text{MgH}_2$  is one of the most promising materials for hydrogen storage due to its high gravimetric and volumetric capacities (7.65 wt.%, 110 g. $\text{H}_2$ /l) and low cost. However, its utilization has been prevented due to its high dissociation temperature, which is much higher than the requirement of the US Department of Energy, and the slow reaction kinetics of this hydride [10–12]. Therefore, different attempts were made in order to improve its adsorption/desorption properties [13–27]. Theoretically, using first principles calculations, it was found that there is a reduction in stability and temperature of decomposition when Mg is combined with several additives [13–16]. By using Kinetic Monte Carlo simulations, it has been shown that  $\text{MgH}_2$  exhibits slow hydrogenation and dehydrogenation kinetics [14]. The addition of small quantities of aluminum and transition metals like Ni, Fe and V can accelerate such kinetics without considerably reducing the associated capacities [14]. Moreover, the single and double substitution of  $\text{MgH}_2$  with lightweight elements, namely boron and lithium, show an improvement in the hydrogen storage properties of the substituted  $\text{MgH}_2$  with a remarkable increase in its gravimetric and volumetric capacities, which exceed those of the pure one [16]. Recently, the effect of magnesium vacancies in the magnesium hydride ( $\text{MgH}_2$ ) has been investigated. The results reveal that the magnesium vacancies and hydrogen doping have a beneficial effect on the hydrogen storage properties of the hydride by decreasing its desorption temperature and stability and increasing its capacities [28,29].

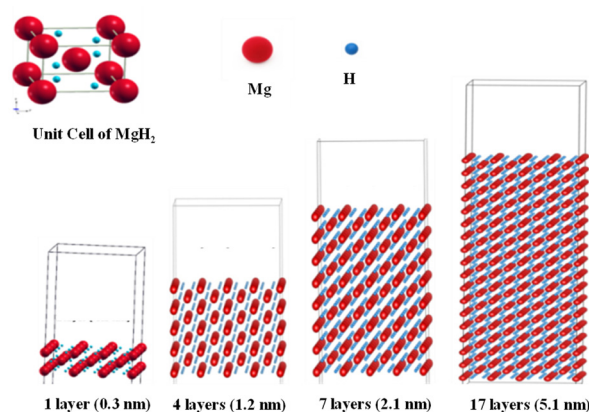
Today, it is well known that nano-structuring materials, with or without straining, lead to new physical phenomena and/or significant improvement of their physical properties for new and existing applications, ranging from the energy field to the bioelectronic and electronic communication fields [30–45]. The straining of nano-sized silicon (Si), for example, increased the plasma wave nonlinearities inside the channel of a traditional Si-based modulation-doped field-effect transistor, which enhanced the coupling and affinity between the Si channel and the electromagnetic radiation in the Terahertz domain [45]. The same has also been shown for compound semiconductors, such as InGaAs, InP and InAs semiconductors. Another example is the nano-structuring of graphene,  $\text{MoS}_2$  and other 2D materials in nanotubes, nanoribbons or quantum dots, which changes the bandgap form and the electronic and optical properties, enabling the use of such materials in many fields, such as biosensing applications (by offering better binding conditions between the nano-structure and the investigated molecules) or Terahertz applications (by providing a tuning dependency between the sizes of the structure and the frequency of operation) [40].

Experimentally,  $\text{MgH}_2$ -based thin films reportedly showed improved hydrogen desorption performance compared to those of bulk  $\text{MgH}_2$  [17–27]. Lu H-B et al. [17] reported that the  $\Delta H$  of hydrogen absorption and desorption for the bulk Mg/ $\text{MgH}_2$  film were significantly high compared to those of the previously prepared 2D nano-structured Mg/ $\text{MgH}_2$  film, with values reaching  $73.9 \pm 0.7$  and  $77.7 \pm 0.8$  kJ/mol, respectively. Barawi et al. [18] and Baldi et al. [19] found that the interface effects between the substrate and the Mg film cause the thickness to be an important factor in the hydrogen absorption kinetics of Mg films. Mooij [20] deduced that the energy of the interface can cause destabilization of the phase of  $\text{MgH}_2$  at the nanoscale when there is a positive difference in interface energy between  $\text{MgH}_2$  and Mg, and causes stabilization in the case of a negative difference. Tao et al. [21] confirmed these findings using the DFT calculations on the dehydrogenation properties of  $\text{MgH}_2/\text{TiH}_2$  multilayers. Moreover, studies conducted over approximately a decade [22,23] showed that 50–100 nm thick  $\text{MgH}_2$ -based thin films have reduced stability. Therefore, the enhancement of the dehydrogenation properties of Mg/ $\text{MgH}_2$ -based nanocomposites undoubtedly depends on the interface misfit and strain energy [24–27]. However, the origins of the kinetic properties and tuned thermodynamics of  $\text{MgH}_2$  of different sizes and under strain had not been thoroughly investigated before. In this paper, the properties of the nano-structured  $\text{MgH}_2$  thin film and the effect of biaxial strain are investigated using first principles calculations.

## 2. Computational Methods

In the present work, DFT-based calculations were performed using the QUANTUM ESPRESSO (QE) code [46]. The generalized-gradient approximation—Perdew–Burke–Ernzerhof (GGA-PBE) functional [47] was used, with a plane wave pseudopotential method available in the QE code. The energy cutoff of 40Ry and the k-point meshes for the Brillouin zone integration were cited as being  $12 \times 12 \times 12$  and  $12 \times 12 \times 1$  for the bulk and thin film of the  $\text{MgH}_2$  structure, respectively. Furthermore, the use of the self-consistent criteria of the density and the energy was conducted with  $10^{-6}$  Ry and  $10^{-8}$  Ry of precision, respectively.

In the first step, the structural parameters of the  $\alpha\text{-MgH}_2$  structure were taken from the experimental results, where  $\alpha\text{-MgH}_2$  crystallized in a rutile structure (space group P42/mnm, No.136) with lattice constants  $a = b = 4.501 \text{ \AA}$  and  $c = 3.010 \text{ \AA}$ , where 2a (0.0, 0.0, 0.0) and 4f (0.304, 0.304, 0.0) sites were occupied by Mg and H atoms, respectively [48]. Then, the structure was fully relaxed with a force convergence of  $10^{-4}$  Ry/Bohr and an energy convergence of  $10^{-7}$  Ry. Once the unstrained unit cell of the bulk  $\alpha\text{-MgH}_2$  was entirely relaxed, the magnesium hydride thin films were modeled by creating  $1 \times 1 \times z$  supercells of different sizes along the Z direction, ranging from four to seventy-six unit cells. For example, a thin film of a thickness of 22 nm was presented with a  $1 \times 1 \times 76$  supercell composed of 308 atoms of magnesium and 154 atoms of hydrogen, while a supercell of  $1 \times 1 \times 4$  represented a thin film of 1.2 nm. Periodic boundary conditions were applied along the X and Y directions in order to make these systems infinite in the (X, Y) plane. In addition, a vacuum of 35 Å was used along the Z-direction to isolate the system of layers, as shown in Figure 1.



**Figure 1.** Structures of bulk and different sizes of  $\text{MgH}_2$  thin films. (The red spheres are magnesium and the blue ones are hydrogen).

The mechanical stress effect on the storage properties of thin film Mg was evaluated by imposing a mechanical biaxial strain along the [100] and [010] directions on the relaxed  $\text{MgH}_2$  (22 nm) thin film according to the following Equation (1) [49–51]:

$$\varepsilon_{xx}(\%) = \varepsilon_{yy}(\%) = \frac{a(b) - a_0(b_0)}{a_0(b_0)} \quad (1)$$

where the lattice constants  $a(b)$  of the  $\text{MgH}_2$  (22 nm) supercell are restricted to different values ranging from  $-2.7\%$  to  $1.8\%$  with a step of  $0.1\%$ , and  $a_0(b_0)$  are the equilibrium lattice constants. The “c” parameter is obtained by relaxing atomic positions under each strain. The negative and positive values of  $\varepsilon_{xx}$  ( $\varepsilon_{yy}$ ) indicate the compression and tensile strain in the  $x$  and  $y$  axes, respectively.

### 3. Results and Discussion

#### 3.1. Size-Dependent Thermodynamic Properties

Before studying the properties of thin film  $\text{MgH}_2$ , the results of the formation enthalpy and desorption temperature of the free strain  $\text{MgH}_2$  unit cell are presented. From the relaxation calculations, the lattice parameters for the bulk structure are  $a = b = 4.520 \text{ \AA}$  and  $c = 3.010 \text{ \AA}$ , which are close to the reported theoretical values [13–16,49] and the experimental ones [48]. The heat of formation for this material was calculated from Equation (2):

$$\Delta H = E_{\text{tot}}(\text{MgH}_2) - E_{\text{tot}}(\text{Mg}) - E_{\text{tot}}(\text{H}_2) \quad (2)$$

where  $E_{\text{tot}}(\text{MgH}_2)$ ,  $E_{\text{tot}}(\text{Mg})$  and  $E_{\text{tot}}(\text{H}_2)$  are the total energies of magnesium hydride, magnesium and molecular hydrogen, respectively. The total energies of the various components were first calculated by using the Plane-Wave Self-Consistent Field (PWSCF) method, and then the heat of formation was deduced using Equation (2).

The desorption temperature can be estimated from Equation (3):

$$T_{\text{des}} = -\frac{\Delta H}{\Delta S} \quad (3)$$

where the entropy is considered to be equivalent to  $\Delta S \approx \Delta S(\text{H}_2) = 130.7 \text{ J/mol.K}$ , which corresponds to the entropy change due to a mole of gas changing to the solid form [52,53]. The different results and energy values are listed in Table 1.

**Table 1.** Energies, formation heat and desorption temperature.

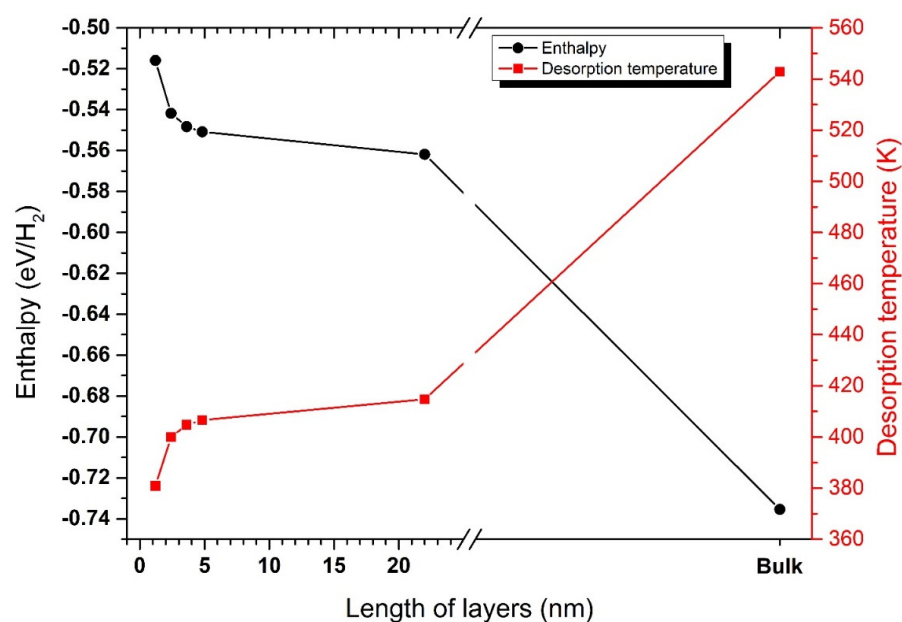
	Energies (eV)	$\Delta H$ (eV/ $\text{H}_2$ )	$T_{\text{des}}$ (K)
Mg (hcp)	−3.07	—	—
$\text{H}_2$	−6.60	—	—
Mg (22 nm)	−235.84	—	—
$\text{MgH}_2$ (bulk)	−17.74	−0.74	542.87
$\text{MgH}_2$ (22 nm)	−1362.37	−0.56	414.75

The calculated enthalpy of formation and desorption temperature of the bulk  $\text{MgH}_2$  are  $-0.74 \text{ eV/H}_2$  and  $542 \text{ K}$ , respectively, which is in good agreement with previous DFT calculations and experimental results [17,54–56]. The magnesium hydride provides a better stability and a high desorption temperature that are much higher than US-DOE requirements [17]. To improve the hydrogen storage properties of this compound, the nano-structuring effect on the thermodynamic properties and hydrogen storage capacities of the  $\text{MgH}_2$  structure was investigated by studying the size dependent hydrogen desorption in Mg hydride thin films with different thicknesses:  $1 \text{ nm}$ ,  $4 \text{ nm}$ ,  $5 \text{ nm}$  and  $22 \text{ nm}$  (Figure 1). To do so, the total energies of Mg thin films, hydrogen molecules and  $\text{MgH}_2$  thin films are calculated first (see Table 1). Subsequently, the heat formation can be deduced using Equation (4) for the thin film model and the corresponding desorption temperature by using Equation (3):

$$\Delta H = E_{\text{tot}}(\text{MgH}_2)_{\text{nano}} - E_{\text{tot}}(\text{Mg})_{\text{nano}} - \frac{n}{2} E_{\text{tot}}(\text{H}_2) \quad (4)$$

where  $n$  is the number of H atoms contained in the  $\text{MgH}_2$  thin film. The effect of the size of  $\text{MgH}_2$  on the desorption temperatures and enthalpy of formation are shown in Figure 2.

By reducing the size, a significant enhancement in the heat of formation is observed from  $-0.74 \text{ eV/H}_2$  for the bulk to  $-0.515 \text{ eV/H}_2$  for the thin film of  $1\text{-nm}$  thickness. This improvement is accompanied by a significant reduction of desorption temperature from  $542 \text{ K}$  to  $380 \text{ K}$ . The last value is in the optimum range  $289\text{--}393 \text{ K}$  for the practical use of fuel cell vehicles [15].



**Figure 2.** Variation of desorption temperatures and enthalpy vs. thickness of MgH<sub>2</sub>.

These results are in line with the reported findings [18,21,55,57], which confirm that the size change on MgH<sub>2</sub> films improves the thermodynamic properties of this compound, owing to the surface properties and the thickness of the 2D nano-structured film. On the other hand, the stability of the MgH<sub>2</sub> thin film was improved compared to other reported works [57–59] by reducing the thickness of the thin film to 1 nm. Additionally, to describe the effect of size on the hydrogen storage capacities of MgH<sub>2</sub>, the gravimetric and volumetric capacities for bulk and different sizes of MgH<sub>2</sub> thin film are calculated and presented in Figure 3. Firstly, the storage capacities of MgH<sub>2</sub> bulk is in good agreement with the values of another study [55]. The obtained results show that MgH<sub>2</sub> has high gravimetric 7.66 wt.% and volumetric 110.14 g.H<sub>2</sub>/l capacities. Secondly, the results show that the volumetric capacity increases from bulk to thin film thickness, whereas the gravimetric capacity remains the same, which indicates that the nano-structuration has a positive effect on the storage capacities.

Assuming that the desorption behavior of the thin films is along the Z axis, it is necessary to aim for nanometer sizes for thin layers in order to have an efficient hydrogen storage medium. The remaining axes (X, Y) were reserved for the application of stress on MgH<sub>2</sub> thin films of 22 nm thickness in order to investigate the effect of such stress on de/hydrogenation properties, as described in the next section.

### 3.2. Strained MgH<sub>2</sub> Thin Film Properties

By reducing the size, we observed a significant enhancement in both the heat of formation and the desorption temperature from  $-0.74$  eV/H<sub>2</sub> and 542 K, respectively, for the bulk to  $-0.515$  eV/H<sub>2</sub> and 380 K, respectively, for the thin film of 1 nm thickness. These results are in good agreement with the reported findings of previously published works [18,55,57], which confirm that the nano-structuring applied in this work by changing the size of MgH<sub>2</sub> films improves the thermodynamic properties of this compound.

In order to simulate the extrinsic effect of the stresses, which can be observed experimentally by the mismatch, a biaxial tensile/compressive strain is applied to the 22 nm thin film along the [100] and [010] directions according to Equation (1). In this case, the lattice constants “a” and “b” of the 22 nm thin supercell are constrained to strain values ranging from  $4.4$  Å to  $4.6$  Å, which correspond to a strain varying from  $-2.7\%$  to  $1.8\%$ , as shown in Figures 4–6.



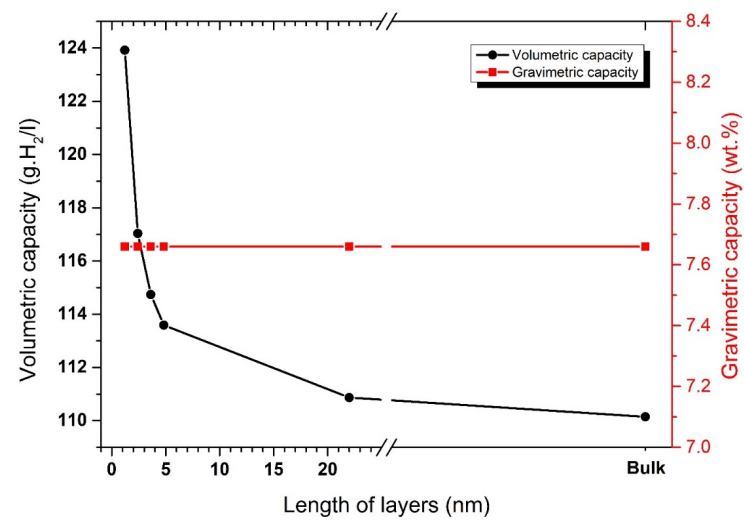


Figure 3. Variation of volumetric and gravimetric capacities vs. thickness of MgH<sub>2</sub>.

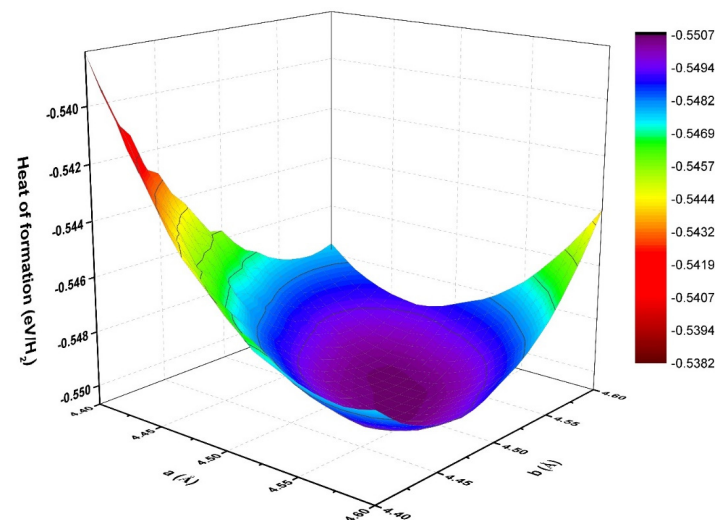


Figure 4. The effect of stress (variation of lattices “a” and “b”) on the stability of MgH<sub>2</sub> thin film (thickness: 22 nm).

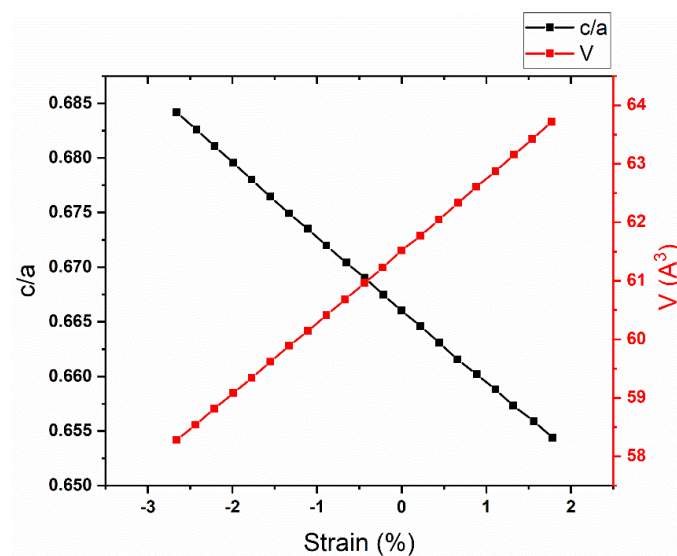
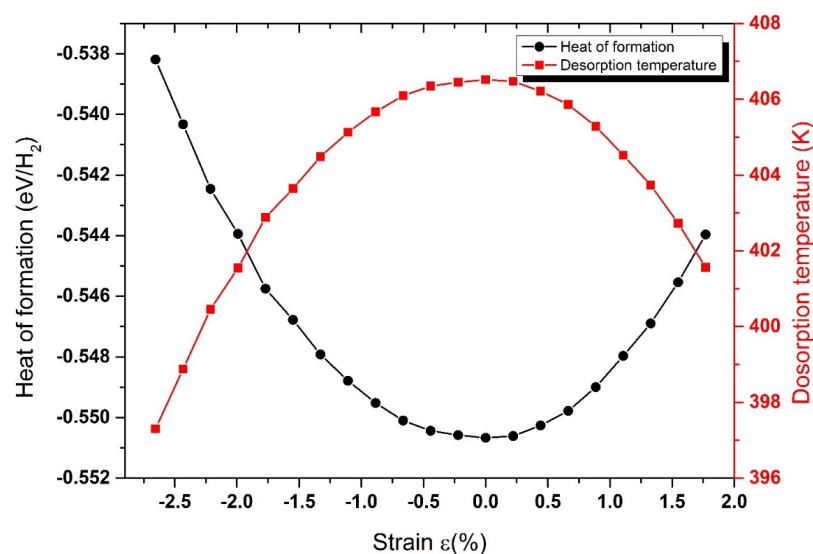


Figure 5. The ratio c/a and supercell volume V as a function of biaxial strain  $\epsilon$  in MgH<sub>2</sub> thin film.



**Figure 6.** Formation enthalpies and desorption temperatures as a function of strain  $\epsilon$ .

Figure 4 shows the stability of a  $\text{MgH}_2$  (22 nm) thin film system under strain. The maximum stability is found for the relaxed parameters, and this stability is reduced considerably when we apply the strain. Moreover, the lattice constant “c” increases under biaxial compressive strain, while it decreases under the tensile one. In other words, under biaxial compressive strain, the lattice constant of the  $\text{MgH}_2$  thin film along the z-axis direction is elongated, which results in an increasing ratio of c/a relative to that (0.667) of the strain-free state, as shown in Figure 5. The maximal ratio of c/a reaches 0.684 at the compressive strain of  $-2.7\%$ . Comparatively, the lattice along the z-axis direction of the nano-structured  $\text{MgH}_2$  shrinks under biaxial tensile strain, yielding a decreasing ratio of c/a relative to that of the strain-free state. A minimum of 0.654 is obtained at the tensile strain of  $+1.8\%$ . Meanwhile,  $\text{MgH}_2$  expands or contracts under biaxial tensile or compressive strain and its cell volume increases or decreases nearly linearly compared to the magnitude of strain as shown in Figure 5. Therefore, it can be derived that either biaxial tensile or compressive strain is likely to cause structural deformation of the  $\text{MgH}_2$  crystal, or its lattice distortion becomes severe with the increasing magnitude of biaxial strain.

The effect of strain on the thermodynamic properties of  $\text{MgH}_2$  thin films (with a thickness of 22 nm) is summarized in Figure 6. This figure shows that the application of strain on  $\text{MgH}_2$  thin film systems reduces their stability compared to that of strain-free ones, which improves their desorption temperatures.

This can be understood as follows: in the case of compressive strain, the distance between atoms is too small; this can lead to a strong interaction and therefore reduces the stability of the system. On the other hand, the H atoms can diffuse easily when we apply a tensile strain, and this consequently reduces the stability of the  $\text{MgH}_2$  thin film with 22 nm of thickness.

The obtained results are partially comparable to the ones found by J. Zhang [49], who states that the strain lowers the decomposition temperature and the stability of the bulk  $\text{MgH}_2$ . Furthermore, the reported variation of formation enthalpy and decomposition temperature observed in the case of compressive strain are totally different than those of the tensile one, whereas in our work, both strains (compressive and tensile) symmetrically increase the enthalpy formation and decrease the desorption temperature for strained  $\text{MgH}_2$  (22 nm) thin film relative to the free strain. The finding’s behaviors are similar to the results of H. Benzidi et al. [60], who showed that  $\text{LiBH}_4$  free of strain leads to a high stability, while compressive or biaxial tensile strain lowers the stability of the  $\text{LiBH}_4$  structure. Finally, due to the contribution of strain, the stability and the decomposition temperature for strained  $\text{MgH}_2$  (22 nm) thin film are reduced compared to those of the strain-free system, which is a great enhancement of  $\text{MgH}_2$  storage properties. The calculation results of this

study show that  $T_{des}$  decreases significantly when increasing the magnitude of strain to a value equal to 398 K, which is close to the optimum range 289–393 K for the practical use of proton-exchange membrane fuel cells (PEMFC) [57].

In conclusion, the nano-structuring of the bulk and the addition of biaxial strain resulted in a beneficial improvement to the storage properties of  $MgH_2$ .

#### 4. Conclusions

In the present paper, the hydrogen storage performances of  $MgH_2$  thin films were studied through the investigation of the effect of size and strain on the thermodynamic properties using first principles DFT calculations implemented in the Quantum-espresso package. Firstly, the studied thermodynamic properties of the bulk  $MgH_2$  present a high stability of  $-0.74$  eV/ $H_2$  and a high decomposition temperature of 542.87 K, which is in good agreement with other theoretical and experimental results. Secondly, the nano-structured  $MgH_2$  thin film improves the thermodynamic properties of the bulk  $MgH_2$  considerably when the decomposition temperature and the stability of the  $MgH_2$  thin film can be reduced to 380 K and  $-0.515$  eV/ $H_2$ , respectively, by reducing its thickness to 1 nm. Finally, in addition to the effect of size, the hydrogen storage properties of  $MgH_2$  (22 nm) thin film were improved by applying a biaxial strain on this system, which caused the stability and the decomposition temperature of the  $MgH_2$  (22 nm) thin film to be reduced from  $-0.56$  eV/ $H_2$  and 414.75 K to  $-0.539$  eV/ $H_2$  and 398 K, respectively, before and after strain. This is a great enhancement of  $MgH_2$  storage properties, even though the strain effect on the thermodynamic properties of the  $MgH_2$  thin film has not been studied previously.

**Author Contributions:** A.E.M. and O.M. conceived/planned the simulations and supervised the findings of this work. A.A., G.T. and M.L. contributed to the configuration of the models. M.L., E.S., A.B., A.K., A.E.K. and H.E.-Z. verified the results and contributed to the interpretation of results. M.L., A.E.M. and O.M. wrote the manuscript with insights from all the authors. A.E.M. secured the funds for this research. All authors discussed the results and contributed to the final manuscript. All authors have read and agreed to the published version of the manuscript.

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