



Article Study on the Homogeneity of Large-Size Blade Lithium-Ion Batteries Based on Thermoelectric Coupling Model Simulation

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Abstract: To improve the energy density of lithium-ion battery packs, lithium-ion batteries are gradually advancing towards large-size structures, which has become one of the dominant development trends in the battery industry. With large-size blade lithium-ion batteries as the research object, this paper develops a high-precision electro-thermal coupling model based on the relevant parameters obtained through basic performance experiments, explores the mechanism of battery inhomogeneity from a simulation perspective, and further proposes a design management method. First of all, the optimal intervals of capacity and temperature, as well as the characteristics of the inhomogeneity distribution for large-size cells, are determined by essential performance and inhomogeneity tests; subsequently, the electrochemical and thermal characteristics of the large-size battery are described precisely through a 3D thermoelectric coupling mechanism model, and the inhomogeneity of the temperature distribution is obtained through simulation; eventually, the optimized cell connection method and thermal management strategy are proposed based on the validated model. As indicated by the findings, the above solutions effectively ease the inhomogeneity of large-size cells and significantly boost the performance of large-size cells under different operating conditions.

Keywords: lithium-ion battery; inhomogeneity; blade battery; thermoelectric coupling model; structural optimization

1. Introduction

Characterized by high energy and power densities, lithium-ion batteries are the workhorse currently among new energy vehicles [1,2]. In limited use spaces, such as electric vehicles, electric boats, and energy storage systems, batteries of higher energy density would be preferred [3]. To cater for this need, the energy density of lithium-ion batteries has to be increased. Researchers in the battery field have sought to increase the energy of the battery by developing higher specific energy cell materials [4], such as high nickel cathode materials [5], lithium metal cathode materials [6], silicon-carbon cathode materials [7], and solid electrolytes, as well as by optimizing the cell and pack structures to increase the energy density of monoblocks and battery packs. In recent years, Panasonic has introduced 21,700 and 4680 cylindrical batteries, with a gradual shift in cell diameter and length to larger cell sizes. Due to the peculiarities of large-size cell structures, cells tend to have inhomogeneous properties in the large-size order, which severely impacts the power performance and durability of the batteries. Therefore, it is essential to understand the inhomogeneity mechanism and accurately model it for the state estimation and management of Li-ion batteries.

Inhomogeneity can be found universally in lithium-ion batteries. It is generally present between solid phase particles [8], in the direction of electrode thickness [9] and between



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Copyright: © 2022 by the authors. Licensee MDPI, Basel, Switzerland. This article is an open access article distributed under the terms and conditions of the Creative Commons Attribution (CC BY) license (https:// creativecommons.org/licenses/by/ 4.0/). the electrode and electrode, triggering inhomogeneous degradation within the battery after a few cycles. Zhang et al. [10] observed that the current distribution is inhomogeneous along the length of the cell, resulting in an inhomogeneous concentration distribution of lithium-ions in phases and inhomogeneous potential distribution [11]. Mühlbaue et al. [12] found that the state of charge (SOC) on the negative surface could vary by up to 10% at larger charging rates. Osswald et al. [13] experimentally demonstrated that the difference in negative electrode potential was up to 515 mV under high magnification discharge conditions, which seriously affected the durability of the battery. The internal and ambient temperature fields of the battery may also exhibit inhomogeneities, due to the inconsistency between internal heat generation and external heat dissipation [14]. As a result of the structure-specific nature of large-size batteries, inhomogeneities in current density, SOC, and temperature, as described above [15], will further affect the homogeneity of mechanical deformation under charge and discharge conditions of large-size Li-ion batteries [16] and lead to the problem of inhomogeneity in the pressure fields within the electrode surface [17]. In general, during the charging and discharging process, physical fields, including pressure, current density, and temperature, will be distributed inhomogeneously. Additionally, this will result in an inhomogeneous distribution of solid-phase and liquid-phase lithium-ion concentrations and potentials, which affects the battery power performance. As the size of the cell structure increases further, the inhomogeneous distribution of multiple physical fields will have a more significant impact on the battery power [18,19].

The durability of the battery is further affected by the uneven degradation within the battery [20]. The rate of the aging decay rate of the battery is closely related to the charge/discharge rate, the internal temperature of the cell, the SOC operating range, and the internal mechanical pressure of the cell [18]. Mussa et al. [21] found that the aging of the positive and negative electrode materials at different locations within the cell was inconsistent through disassembly of the shell battery, with the aging of the electrode materials at the bends being the most severe. The P2D model is one of the most prominent ways to model the mechanism of the lithium-ion battery, which can accurately realize the state of charge and lithium-ion distribution and kinetic processes in lithium-ion batteries, and scholars have conducted a series of studies on this subject [1-3]. Based on the results of model simulation, Arunachala et al. [22] developed an electrochemical model and suggested that the current and temperature inhomogeneity would cause local abuse, which, as a result, would worsen the inhomogeneity of the degree of aging within the battery and could eventually cause localized abuse failure over a large area, accelerating the aging process of the battery. Cannarella et al. [23,24] observed, in locally closed-hole regions fabricated on diaphragms, that the lithium precipitation area of the negative graphite electrode increased with cycling due to local concentration of current density. Some scholars constructed two-dimensional models considering the failure behavior of local abuse caused by mechanical–electrochemical–thermal inhomogeneities. Erhard et al. [13] developed a battery thermoelectric coupling model taking into account the spatial distribution to realize the simulation of potential, temperature, and pressure distribution characteristics inside the battery under different operating conditions and identified inhomogeneities of these quantities of states in the spatial distribution. In contrast, Zhu et al. [25] established a simulation model of localized lithium precipitation abuse in batteries triggered by localized temperature hot spots, but the accuracy and robustness of this model lacks experimental verification, due to the absence of intuitive data on the performance of large-size batteries in actual use (that is, the capacity, internal resistance, entropy heat coefficient, specific heat capacity, and thermal conductivity).

In summary, inhomogeneous properties induced by multi-physical fields and different battery structures lead to inhomogenous aging within the battery, even triggering local abuse failure and accelerating the aging process. This problem will become more pronounced as the battery size grows. It emphasizes the necessity to analyze the nonhomogeneous characteristics of batteries from multiple perspectives, including experimental tests, mechanism research, and model building. In this paper, conducting basic tests on large commercial batteries, such as the single cell capacity and entropy thermal coefficient, the performance of non-homogeneity of large-size batteries under constant current charge and discharge conditions is highlighted. Meanwhile, the non-homogeneity temperature distribution of large-size batteries is also analyzed and summarized. Based on the pseudo-two-dimensional (P2D) model, a three-dimensional thermoelectric coupling model is constructed, and its accuracy is verified at various temperatures and charge/discharge rates. Based on the verified three-dimensional thermoelectric coupling model, the mechanism of large-size battery inhomogeneity causation and the structural optimization and management methods to alleviate large-size battery inhomogeneity are proposed, and a comparison of the inhomogeneity distribution of the cells before and after optimization is made.

2. Theoretical Analysis

2.1. Pseudo Two-Dimensional (P2D) Model

The mechanistic model used in this paper is built based on the pseudo-two-dimensional (P2D) model developed by Doyle and Newman et al. [26,27]. The P2D model can accurately realize the lithium-ion distribution states and kinetic processes in Li-ion batteries, thus accurately simulating the electrochemical characteristics of Li-ion batteries under different temperatures and C-rate charging. The electrochemical control equations are as follows.

The detailed charge conservation equation in solid and liquid phases is given below. The charge balance equation in the solid phase follows Ohm's law, as shown in Equation (1). In this equation, σ_s^{eff} is the effective conductivity of the solid phase, ϕ_s denotes the potential of the solid phase, a_s represents the specific interface area of the particle, F is the Faraday constant, and j denotes the current density. The charge balance equation in the liquid phase is demonstrated in Equation (2). In this equation, κ_e^{eff} represents the effective conductivity of the electrolyte, ϕ_e is the electrolyte potential, R denotes the gas constant, T is the cell temperature, t_+ is the number of lithium-ions transferred, and c_e is the lithium-ion concentration in the electrolyte.

$$\frac{\partial}{\partial x} \left(\sigma_s^{eff} \frac{\partial}{\partial x} \phi_s \right) - a_s F j = 0 \tag{1}$$

$$\frac{\partial}{\partial x} \left(\kappa_e^{eff} \frac{\partial}{\partial x} \phi_e \right) + \frac{2RT(t_+ - 1)}{F} \frac{\partial}{\partial x} \left(\kappa_e^{eff} \frac{\partial}{\partial x} \ln c_e \right) + a_s F_n = 0$$
(2)

In the solid-liquid two-phase matter conservation equation, the matter conservation of lithium-ions in the active electrode material is described by Fick's law, as indicated in Equation (3). C_s refers to the solid-phase concentration of lithium-ions in the particles, D_s denotes the solid-phase diffusion coefficient, and r represents the particle radius. The lithium-ion concentration distribution in the electrolyte is presented in Equation (4). In this equation, ε_e is the volume fraction of the electrolyte, and D_e^{eff} refers to the effective diffusion coefficient in the electrolyte.

$$\frac{\partial c_{\rm s}}{\partial t} - \frac{D_{\rm s}}{r^2} \frac{\partial}{\partial r} \left(r^2 \frac{\partial c_{\rm s}}{\partial r} \right) = 0 \tag{3}$$

$$\frac{\partial}{\partial t}(\varepsilon_{\rm e}c_{\rm e}) = \frac{\partial}{\partial x} \left(D_{\rm e}^{\rm eff} \frac{\partial}{\partial x} c_{\rm c} \right) + (1 - t_{+}) a_{\rm s} j \tag{4}$$

In the solid-liquid two-phase lithium-ion concentration distribution equation, the distribution of lithium-ion concentration diffusion process in both solid-liquid phases satisfies Fick's second law, as shown in Equations (5) and (6). In these two equations, D_s is

the solid-phase diffusion coefficient of lithium-ions, ε_e represents the volume fraction of electrolyte, and D_e^{eff} refers to the effective diffusion coefficient of the liquid phase.

$$\frac{\partial c_s}{\partial t} - \frac{D_s}{r^2} \frac{\partial}{\partial r} \left(r^2 \frac{\partial c_s}{\partial r} \right) = 0 \tag{5}$$

$$\frac{\partial}{\partial t}(\varepsilon_e c_e) = \frac{\partial}{\partial x} \left(D_e^{eef} \frac{\partial}{\partial x} c_e \right) + (1 - t_+) a_s j_n \tag{6}$$

In the kinetic equations for the electrochemical reaction at the two-phase interface, for the intercalation reaction of lithium at the cathode and anode, the local current density in the active material region satisfies the Bulter–Volmer (BV) equation, as can be seen in Equation (7). In the equation, i_0 refers to the current exchange density, α_a and α_c stand for the anodic transfer coefficient and cathodic transfer coefficient, respectively, and η represents the reaction overpotential.

$$j_n = i_0 \cdot \left[\exp\left(\frac{\alpha_a F \eta}{RT}\right) - \exp\left(\frac{\alpha_c F \eta}{RT}\right) \right]$$
(7)

2.2. Li-Ion Battery Heat Production and Heat Transfer Model

It can be observed from the P2D model that the equilibrium potential of the battery is strongly influenced by temperature. The charging and discharging process of the battery is accompanied by heat absorption and exothermic phenomena, which affect the performance of the battery potential. Therefore, to establish a more accurate mechanism model to characterize Li-ion batteries under different operating conditions, further coupling of the heat generation and heat transfer equations of the battery is needed.

In the lithium-ion battery heat production equation, according to the literature [28–31] on the thermal model of lithium-ion batteries, the lithium-ion battery heat production Q is mainly divided into two types, namely reversible heat Q_{rev} and irreversible heat Q_{irrev} , as shown in Equation (8). Reversible heat Q_{rev} is the energy absorbed or released by the lithium intercalation material during lithium intercalation and deintercalation. Its value is proportional to the entropy heat coefficient of the electrode, which is also known as the reaction heat Q_{rea} , as demonstrated in Equation (9). The irreversible heat, on the other hand, is divided into ohmic heat Q_{ohm} and polarization heat Q_{act} . The ohmic heat Q_{ohm} refers to the heat production when electrons and lithium-ions collide with the transport medium during the internal transport of the battery, as shown in Equation (10).

$$Q = Q_{\rm rev} + Q_{\rm irrev} = Q_{\rm rea} + (Q_{\rm ohm} + Q_{\rm act})$$
(8)

$$Q_{\rm rea} = a_s F j_n T \frac{dU}{dT} \tag{9}$$

$$Q_{ohm} = \sigma_s^{eff} \left(\frac{\partial \phi_s}{\partial x}\right)^2 + \sigma_e^{eff} \left(\frac{\partial \phi_e}{\partial x}\right)^2 + \frac{2\sigma_e^{eff}RT}{F}(1-t_+)\frac{\partial \ln c_e}{\partial x}\frac{\partial \phi_e}{\partial x} + I(t)R_c$$
(10)

 Q_{act} , the heat of polarization, then refers to the heat of the activation energy barrier that needs to be broken to change the equilibrium state of the electrode embedded lithium, as shown in Equation (11). Since Equation (11) contains an overpotential term, the heat of polarization is coupled to the BV equation, and the heat of polarization is related to the potential two-phase distribution and the electrochemical reaction at the two-phase interface.

$$Q_{\rm act} = a_{\rm s} F j_n (\phi_{\rm s} - \phi_{\rm c} - U - F j_n R_{\rm SEI})$$
⁽¹¹⁾

The heat transfer equation of a lithium-ion battery depends not only on the heat generated by itself, but also on the heat dissipation in the environment, as shown in Equation (12), which is the heat transfer equation of the battery, also known as the energy conservation equation. In this equation, ρ represents the density of the battery, C_p denotes

the specific heat capacity of the battery, h is the convective heat transfer coefficient of the battery surface, and T_{ext} refers to the ambient temperature. This equation states that the rate of temperature change of the battery is equal to the difference between the rate of heat production and heat dissipation of the battery.

$$\rho C_p \frac{\partial T}{\partial t} = Q - hA_s(T_{ext} - T)$$
(12)

3. Experimental Setup

In this study, a commercial prismatic power battery with a nominal capacity of 135 Ah was adopted, with LiFPO₄ as the cathode and graphite as the negative electrode. In the subsequent basic performance test, the ambient temperature would have to be adjusted according to operating conditions. The SC-360-CC-3 large-size temperature chamber from SANWOOD Technology was applied as the explosion-proof temperature chamber for battery experiments, and the CE-7004-100V300A-R28GC power battery test system produced by Neware was adopted.

3.1. Basic Performance Testing of Large-Size Lithium-Ion Batteries

A total of 12 thermocouples (type GG-KK-36-SLE manufactured by ETA) were arranged on the front and back of the battery to record temperature changes at different locations, as shown in Figure 1. There were 9 thermocouples on the front and back of the battery, namely 2, 3, 9, and 10 near the positive and negative lugs, 8 and 4 at 1/4 and 3/4 of the length of the battery, 5, 6, and 7 at 1/2 of the length of the battery, and 1 and 11 at the center of the positive and negative battery lugs, as shown in Table 1.



Figure 1. Large-size battery thermoelectric coupling layout diagram.

T °C/C Rate	0.2C	0.3C	0.5C	1C	1.5C	2C
-10	Cha and Dch	Cha	Cha and Dch	Dch	/	/
0	Cha and Dch	Cha	Cha	Dch	/	Dch
25	Cha and Dch	/	Cha and Dch	Cha and Dch	Cha and Dch	Cha and Dch
45	/	/	/	Cha and Dch	Cha and Dch	Cha and Dch

In order to study the effect of temperature and multiplier on the capacity of largesize batteries, the experiments in this section included four temperature and six chargedischarge rate conditions, each consisting of a charge capacity test and a discharge capacity test and all adopting constant current (CC) charging and discharging methods, as shown in Table 1.

Following the capacity test, the standard capacity of the battery at each temperature was tested, and the standard capacity test was measured with the 1/3C CC-CV work step.

3.2. *Experimental Calibration of Thermoelectric Coupling Model Related Parameters* 3.2.1. Equilibrium Potential of Positive and Negative Electrodes and Acquisition of Monomer OCV Curve

In order to obtain an equilibrium potential between the positive and negative electrode materials, it is necessary to fabricate positive-lithium metal button cells and negativelithium metal button cells. Prior to fabrication of the button cells, the electrode sheet and diaphragm were treated as follows. The disassembled positive and negative electrode sheets and fresh PP diaphragms (produced by Clutech) were transferred to an argon environment glove box with a water content of less than 0.1 ppm. The positive and negative plates and the diaphragm on the side with the active substance removal were cut into discs 12 mm and 19 mm in diameter, respectively. To remove lithium salts and impurities from the surface of the electrodes, the punched round electrodes were soaked in a vial of DMC (dimethyl carbonate) solution for 3 to 6 h, which were then removed and dried. The punched round diaphragms were transferred to an oven for 12 h. The treated electrodes and diaphragms were stacked as follows. First of all, the positive/negative electrodes of the large-size cell were placed on top of the positive shell, with an inner diameter of 19 mm, and then the diaphragm was placed on top of the positive electrode. The diaphragm was placed on top of the positive electrode sheet, and 3-5 drops of 1.0 mol LiPF6-EC:EMC: DMC (1:1:1) electrolyte were added. The lithium metal sheet of 16 mm diameter metal sheet-metal spacer-funnel type shrapnel was placed on top of the diaphragm in sequence, and finally, the negative shell was covered and pressed with a press.

The fabricated sets of positive-lithium metal and negative-lithium metal button cells were placed in the button cell test machine at a temperature of 25 °C for more than 12 h to allow the electrolyte to fully infiltrate the active material of the button cell. Afterwards, the positive and negative coin cells were tested with a charge/discharge current of 2 mA (approximately 0.05 C–0.07 C in proportional conversion), and the results of the equilibrium potential of positive and negative materials can be seen shown in Figure 2.



Figure 2. Equilibrium potential curve of large-size cell.

3.2.2. Identification of Electrode Parameters Based on the Dual Water Tank Model

In the subsequently constructed thermoelectric coupling model, the theoretical capacity of the electrode material and the electrode SOC operating interval would be specified, and the cell OCV would be calculated by the model based on the initially set interpolation values, so the above parameters would have to be determined by our previously proposed dual water tank model [32,33].

The capacity of $C_{Fullcell}$ as a single cell corresponds to the test data of the single cell in Section 3.1 and to the test data of positive and negative button cells in Section 3.2.1, respectively, with the charging current *I* set to 1/20C, and the identification results are shown in Figure 3.



Figure 3. Identification results of Large-size battery dual tank model.

The maximum concentration of embedded lithium in the positive and negative materials and the total area of active materials are calculated by Equations (13) and (14), based on the theoretical capacity of the positive and negative materials determined from the dual tank model. In these equations, $Cap_{neg/pos}$ is the theoretical capacity of positive and negative electrodes, 3600 refers to the conversion of capacity unit Ah to charge unit Coulomb (C), F denotes the Faraday constant, $\varepsilon_{s, neg/pos}$ represents the volume fraction of solid phase, and $L_{neg/pos}$ represents the thickness of active material (as indicated by Table 2). As represents the total area of active material, W_{cell} and H_{cell} represents the width of large-size cell, n_{layers} denotes the number of positive electrodes, and since each electrode has active material on the front and back side, the number of electrodes needs to be multiplied by 2 when calculating the total area of active material.

$$c_{s,neg/pos,max} = \frac{3600 \cdot Ca_{neg/pos}}{F\varepsilon_{s,neg/pos}A_s L_{neg/pos}}$$
(13)

$$A_s = 2W_{batt}H_{batt}n_{layers} \tag{14}$$

Symbol	Meaning of Parameters	Unit	Value		
Thermal model parameters					
F	Faraday constant	$C \cdot mol^{-1}$	96,485		
R^0	Ideal gas constant	$J \cdot mol^{-1} \cdot K^{-1}$	8.3145		
T_{ref}	Reference temperature	K	298.15		
ρ_{neg}	Density of negative active material	kg∙m ⁻³	2300		
$ ho_{pos}$	Density of positive active material	kg⋅m ⁻³	3600		
ρ_{sep}	Diaphragm density	kg∙m ^{−3}	1108.98		
$\rho_{neg,cc}$	Fluid density of negative set	kg∙m ^{−3}	8960		
$ ho_{pos,cc}$	Positive electrode fluid density	kg∙m ^{−3}	2700		
ρ_{batt}	Battery cell density	kg∙m ^{−3}	2600		
$k_{T,neg}$	Thermal conductivity of negative active material	$W \cdot m^{-1} \cdot K^{-1}$	1.04		
$k_{T,pos}$	Thermal conductivity of positive electrode active material	$W \cdot m^{-1} \cdot K^{-1}$	1.58		
$k_{T,sep}$	Thermal conductivity of diaphragm	$W \cdot m^{-1} \cdot K^{-1}$	0.344		
k _{T,neg,cc}	Thermal conductivity of negative current collector	$W \cdot m^{-1} \cdot K^{-1}$	400		
$k_{T,pos,cc}$	Thermal conductivity of positive current collector	$W \cdot m^{-1} \cdot K^{-1}$	238		
k _{T,batt}	Thermal conductivity of battery cell	$W \cdot m^{-1} \cdot K^{-1}$	0.98		
$C_{p,neg}$	Specific heat capacity of negative active material	$J \cdot kg^{-1} \cdot K^{-1}$	750		
$C_{p,pos}$	Specific heat capacity of positive electrode active material	$J \cdot kg^{-1} \cdot K^{-1}$	881		
$C_{p,sep}$	Specific heat capacity of diaphragm	$J \cdot kg^{-1} \cdot K^{-1}$	1978.16		
C _{p,neg,cc}	Specific heat capacity of negative current collector	$J \cdot kg^{-1} \cdot K^{-1}$	385		
$C_{p,pos,cc}$	Specific heat capacity of positive current collector	$J \cdot kg^{-1} \cdot K^{-1}$	900		
$C_{p,batt}$	Specific heat capacity of battery cell	$J \cdot kg^{-1} \cdot K^{-1}$	983.9		
h _{25°C}	Heat transfer coefficient at 25 °C	$W \cdot m^{-2} \cdot K^{-1}$	12		
$h_{45^{\circ}C}$	Heat transfer coefficient at 45 $^\circ C$	$W \cdot m^{-2} \cdot K^{-1}$	16		
$h_{0^{\circ}C}$	Heat transfer coefficient at 0 °C	$W \cdot m^{-2} \cdot K^{-1}$	8		
<i>h</i> ₋10°C	Heat transfer coefficient at -10 $^\circ C$	$W \cdot m^{-2} \cdot K^{-1}$	10		

Table 2. Model parameters of electrochemical-thermal coupling mechanism and thermal physicalparameters of large-size cells.

Electrochemical model parameters

Symbol	Meaning of parameters	Unit	Value		
0 9 11 2 01	including of parameters		Neg	Sep	Neg
L	Thickness	m	65.5×10^{-6}	$30 imes 10^{-6}$	72.5×10^{-6}
R_s	Particle radius	m	$30 imes 10^{-6}$	/	$10 imes 10^{-6}$
\mathcal{E}_S	Active material volume fraction	1	0.65	/	0.55
ε_e	Electrolyte volume fraction	1	0.3	0.4	0.2
a_s	Specific surface area	m^{-1}	$3\varepsilon_{s,neg}/R_{s,neg}$	/	$3\varepsilon_{s,pos}/R_{s,pos}$
A_s	Battery area	m ²	5.6618		
P_c	Battery circumference	m	1.971		
Lneg,cc	Thickness of negative current collector	m	$10 imes 10^{-6}$		
Lneg,cc	Positive current collector thickness	m	$15 imes 10^{-6}$		
L _{batt}	Battery cell thickness	m	$1.93 imes10^{-4}$		
W _{batt}	Battery cell width	m	$8.25 imes 10^{-2}$		
H_{batt}	Battery cell length	m	$9.03 imes10^{-1}$		
W _{tab}	Polar ear width	m	$7.15 imes 10^{-2}$		
H_{tab}	Polar ear length	m	$2 imes 10^{-2}$		
<i>n_{layers}</i>	Number of poles	1	38		
$\dot{x_0}$	Positive Stoichiometric Ratio of cell at 0% SOC	1		0.0073	
x_{100}	Positive Stoichiometric Ratio of cell at 100% SOC	1		0.9501	
y_0	The negative stoichiometric ratio of the cell at 0% SOC	1		0.9329	
Y 100	The negative stoichiometric ratio of the cell at 0% SOC	1		0.2904	
σ_s	Solid phase conductivity	S⋅m ⁻¹	100	/	91
σ_s^{eff}	Solid phase effective conductivity	$S \cdot m^{-1}$	$\epsilon_{s,neg}^{1.5}\sigma_{s,neg}$	/	$\varepsilon_{s,pos}^{1.5}\sigma_{s,pos}$
$\sigma_{s,cc}$	Current collector electronic conductivity	$S \cdot m^{-1}$	$1.4 imes10^8$	/	$3.774 imes 10^7$
$\sigma_{s,tab}$	Pole-ear electronic conductivity	S⋅m ⁻¹	$7 imes 10^7$	/	$1.5 imes10^7$
κ _e	Liquid phase conductivity	$S \cdot m^{-1}$	$10^{-8.43-7}$	$\frac{54}{T-5\times 10^{-3}\times c-299}$ -	$2.2 \times 10^{-4} \times c$
κ_e^{eff}	Liquid effective conductivity	$S \cdot m^{-1}$	$\varepsilon_{e,neg}^{1.5}\kappa_e$	$\varepsilon_{e,sep}^{1.5}\kappa_{e}$	$\varepsilon_{e,pos}^{1.5}\kappa_{e}$

Symbol	Meaning of Parameters	Unit	Value	
D_e^{eff}	Effective diffusion coefficient of liquid phase	$\mathrm{m}^2 \cdot \mathrm{s}^{-1}$ $\varepsilon_{e,neg}^{1.5} D_e$	$\varepsilon_{e,sep}^{1.5}D_e$	$\varepsilon_{e,pos}^{1.5}D_e$
$D_{s,ref}$	Solid diffusion coefficient at reference temperature	$m^2 \cdot s^{-1}$ 1 × 10 ⁻¹¹	1	$4 imes 10^{-14}$
E_{a,D_s}	Activation energy of Ds when temperature changes	kJ·mol ^{−1} 358.1	/	8.8
k _{ref}	Reaction rate coefficient at reference temperature	$\mathrm{m\cdot s^{-1}}$ 1 $ imes$ 10 ⁻¹⁰	/	$5 imes 10^{-11}$
$E_{a,k}$	Activation energy of k when temperature changes	$kJ \cdot mol^{-1}$ 2.1	/	5.1
$C_{s,max}$	Maximum lithium-ion concentration of material	mol∙m ⁻³ 28,244	/	23,118
С _{е,0}	Initial concentration of liquid lithium-ion	mol⋅m ⁻³	1200	
<i>x,y</i>	Lithium-ion stoichiometric ratio	1 0.0085, 0.771	/	0.0033, 0.998
t_+	Lithium-ion migration number	1	0.382	
α_a, α_c	Lithium-ion transfer coefficient	1 0.5, 0.5	/	0.5, 0.5
R_{SEI}	SEI film resistance	$\Omega \cdot m^2 = 4 imes 10^{-3}$	/	/
dU/dT	entropy heat coefficient	V Figure 3	/	/
U_{ref}	Equilibrium potential	mV·K ^{-1} Figure 2	/	Figure 2

Table 2. Cont.

3.3. Experiment of Acquiring Parameters of Thermoelectric Coupling Model

3.3.1. Measurement of Basic Parameters of Thermal Physical Properties

The monomer entropy heat factor $\Delta E/\Delta T$ (mv/K) exerts a significant impact on the cell heat generation model and a test of the monomeric heat coefficient of the largesized cell is required. Based on the results of our previous work [34], the obtained cell monomer entropy heat coefficients are demonstrated in Figure 4. To accurately address the coupling between the electrochemical and thermal properties in the mechanism model, it is also essential to understand the relationship between the entropy-heat coefficient of the dU/dT electrode material and SOC of the electrode. Since the measurement of material entropy heat coefficients requires a high level of experimental methodology and accuracy, this paper provides the entropy heat coefficients for graphite and LFP materials based on those obtained from the COMSOL Multiphysics^{®®} (version 5.4) large-size battery materials library.

The diffusion coefficients of electrolyte D_e and electrolyte ion conductivity κ_e are affected by temperature *T* and electrolytic lithium-ion concentration *c* at various points inside the cell. The electrolyte parameters obtained based on references [29,35] and parameters of the material library with suitable modifications are presented in Equations (15) and (16) hereafter.

$$D_e = 10^{-8.43 - \frac{54}{T - 5 \times 10^{-3} \times c - 299} - 2.2 \times 10^{-4} \times c}$$
(15)

$$\kappa_e = \left(-10.5 + 0.074T - 6.96 \times 10^{-5}T^2 + 6.68 \times 10^{-4}c - 1.78 \times 10^{-5}cT + 2.8 \times 10^{-8}cT^2 + 4.94 \times 10^{-7}c^2 - 8.86 \times 10^{-8}c^2T^2\right) \times 10^{-6}c$$
(16)

In the solid heat transfer model, the monomer, specific heat capacity of each part of the material, thermal conductivity, and heat transfer coefficient can significantly affect heat generation at each point of the cell, and the parameters of this part were calibrated according to Refs. [34,36,37] and the experimental results, as detailed in Table 2.

3.3.2. Electrochemical Model Parameters Determination

According to the P2D model, up to twenty-five parameters from the coupled multiphysics field model have diverse and complex sources. Currently, the main categories are cell design parameters, ion transport characteristics parameters, and interface reaction kinetics parameters. In order to improve the accuracy of the model and the modeling cost, the model parameters should be obtained experimentally whenever possible to reduce the number of post-calibrations. In this paper, the parameters are mainly accessed through disassembly measurements, experimental tests, parameter identification, and literature references. Combined with the previous work [34], the relevant parameters of the electrochemical model can be seen in Table 2.



Figure 4. Entropy thermal coefficient test results of large-size battery monomer.

4. Results and Discussion

4.1. Experimental Test Results of Basic Performance of Large-Size Batteries

As shown in Figure 5, the results of the two capacity tests in Section 3.1 are graphically presented by a capacity charge/discharge rate temperature 3D plot. Additionally, the findings can be reached as follows.

(1) The influence of temperature on charging and discharging capacity can be substantial, especially when the temperature falls below 0 degrees, and both the discharging and charging capacity are significantly reduced. In contrast, temperatures above 25 °C, such as the room temperature, exert little influence on the capacity of this large-size battery.

(2) Generally, the optimum range of temperature for battery capacity performance is expected to be between 25 °C and 45 °C, while the optimum range of charge/discharge rate is between 0.2 C and 1.5 C.

(3) In addition, it can be found in Figure 5a,b that the charging capacity decreases when the low-temperature charging rate increases (solid black arrows in the figure), while the discharging capacity increases slightly when the low-temperature discharge rate increases in Figure 5c,d (solid blue arrows in the figure). These results indicate that, when designing a low temperature charging strategy for this large-size battery, slow charging at low multiples is preferred and common methods, such as coolant heating, heat pipe heating, and AC heating, to increase the battery temperature should be adopted when large-size batteries have to work in low temperature conditions, as this will help the battery to perform at a higher range.



Figure 5. Capacity Charge/discharge Rate Temperature 3D plot of large-size battery for two tests. (**a**,**b**) Charging capacity (first/second test); (**c**,**d**) Discharging capacity (first/second test).

4.2. Model Verification

In order to better understand the inhomogeneity mechanism of the large-size cells, a three-dimensional thermoelectric coupling mechanism model for this large-size cell is developed in this section. It is developed in COMSOL Multiphysics^{®®} software (version 5.4), based on pseudo-two-dimensional theory and solid-state heat transfer theory, with reference to our previous thermoelectric coupling modeling process and the way of obtaining model parameters.

A comparison of CC charging voltage curves for different operating conditions of the 3D electrochemical-thermal coupling model with the experimental results is demonstrated Figure 6. It can be observed that, after adding two cell length and width dimensions to the P2D model, and including the possible influence of the collector conductivity on the external monomer characteristics, the simulated voltage results under different operating conditions still match well with the experimental values. The RMSE (root mean square error) voltages of the different C-rate simulation results at 25 degrees and 45 degrees are lower than 7.5 mV. The models at 0 degrees and -10 degrees show some variability at high SOC, but the overall accuracy is good, indicating that the 3D thermoelectric coupling model can describe the charging characteristics of large-size batteries under different operating conditions.

Figure 7 validates the temperature response of the 3D model for different charging conditions. It should be noted that, for low rate charging, the machine model results in the smallest simulation error, so the maximum charging rate at four temperatures (that is, 25 °C and 45 °C 2 C conditions and 0 °C and -10 °C 0.5 C conditions) are used here as an example for validation, with the temperature validation locations being three temperature points on the midline of the large surface of the large-size battery, i.e., Figure 1 (a) in points 4, 6, and 8. As can be seen from the figure, the temperature results of the 3D model still match the experimental results. Due to the increased heat dissipation at the pole lugs and

the low self-generated heat of the battery at the end of charging, the heat dissipation has a greater impact on the battery temperature, ultimately resulting in a model temperature at the end of charging that is slightly lower than the experimental temperature. However, the average temperature error is less than $2 \,^{\circ}$ C. In addition, there is temperature inhomogeneity in the experimental temperature chamber under low-temperature conditions, with the temperature decreasing sequentially from point 4 to point 6 near the positive pole under low temperatures. However, the trend of temperature rise is the same.



Figure 6. The 3D model charging voltage validation for different operating conditions (shape-experimental data, solid line-model results). (**a**) 25 °C; (**b**) 45 °C; (**c**) 0 °C; (**d**) -10 °C.

To further verify the temperature inhomogeneity of the 3D model, the results of the 3D temperature distribution at the initial moment of charging and the cut-off moment under four charging multipliers are revealed in Figure 8, taking an operating condition of 25 °C as an example. Figure 8 suggests that the battery surface temperature is homogeneous at the initial charging moment, and the temperature inhomogeneity at the cut-off moment is highlighted. With the increase of the charging multiplier, the maximum temperature difference at the charging cut-off moment increases from 0.4 °C at 0.5 C to 8.8 °C at 2 C, which is consistent with the experimental results. The temperature inhomogeneity at the charging cut-off point is higher at the lug side and lower in the middle under conditions greater than 1 C, which is also the same as the inhomogeneity test results in Section 3.1.

In summary, the three-dimensional thermoelectric coupling model of large-size batteries can accurately describe the electrochemical and thermal characteristics of the batteries under different operating conditions and can also accurately simulate the surface temperature inhomogeneities of the large-size batteries, and as such, it has laid the foundation for the subsequent in-depth analysis of the model results.



Figure 7. The 3D model midline temperature validation for different charging conditions. (a) $25 \degree C$ 2 C; (b) $45 \degree C 2 C$; (c) $0 \degree C 0.5C$; (d) $-10 \degree C 0.5 C$.



Figure 8. Results of the three-dimensional temperature distribution model for charging at 25 °C 0.5 C, 1 C, 1.5 C, 2 C. (**a**,**c**,**e**,**g**) charging initial moment; (**b**,**d**,**f**,**h**) charging cutoff moment.

4.3. Mechanism Analysis Based on the Inhomogeneity of Electrothermal Coupling Model

The difference of SOC on the negative separator side at the charging cut-off moment is analyzed at different charging rates of 25 °C as an example, as shown in Figure 9, with the upper part of the two-dimensional diagram being close to the negative terminal lug side and the lower part being close to the positive terminal lug side. It can be seen that there is a difference of about 1% in the SOC distribution on the negative diaphragm side at the charging cutoff moment in the 0.5 C condition, which is slightly higher on the positive lug side and slightly lower on the negative lug side. The SOC distribution at 1 C and 0.5 C on the negative diaphragm side is similar to that at 0.5C, while at 2 C condition, the SOC is higher on the two lug sides and lower in the middle, with a SOC difference of around 7%. The comparison with Figure 8 indicates that this is the same situation as the uneven temperature distribution, where the cell temperature is also higher at the higher SOC location. This is mainly a result of the convergence of the electron currents entering the active material on the lug side, resulting in higher utilization of the electrode on the lug side and more heat generation during the charging process, as well as the heat generation effect of ohmic polarization at the lug position, which together result in a faster temperature rise on the lug side.



Figure 9. Difference in SOC on the negative diaphragm side at the charging cut off moment for different operating conditions at 25 °C. (**a**) 0.5 C; (**b**) 1 C; (**c**) 1.5 C; (**d**) 2 C.

In addition, the intercalation current density distribution of the lithium-ion embedding reaction on the negative diaphragm side at the charging cut-off moment is also demonstrated in Figure 10. It can be seen that the inhomogeneity of the intercalation current density is the opposite of the SOC inhomogeneity and temperature inhomogeneity. As of the moment of charging, the lithium-ion embedding current density is lower on the bipolar ear side, where the electron currents converge due to the higher SOC and the lower lithium-ion embedding reaction rate, while the interlayer current density is relatively high at the location of the protruding surface in the middle, where the SOC is lower. As the multiplicity increases, the current density difference at 0.5C is about 2 A/m², while it reaches 12 A/m² at 2 C. However, a more significant current density difference at 2 C also results from the charging current difference, with a normalized current density difference of approximately 0.57 at 0.5 C working conditions, as opposed to approximately 0.75 at 2 C working conditions.

4.4. Unevenness Management Method Based on Electrothermal Coupling Model

Experimental and simulation results demonstrate that there is significant temperature inhomogeneity in the charging and discharging conditions of large-size batteries. Therefore, an optimization method of battery connection to reduce the temperature inhomogeneity, as well as a thermal management strategy to adjust the heat transfer coefficient at different locations based on a three-dimensional thermoelectric coupling model (based on the 25 °C 2 C charging model), is proposed in this section.



Figure 10. Difference in current density of the negative diaphragm side intercalation at the charging cutoff moment for different operating conditions at 25 °C. (a) 0.5 C; (b) 1 C; (c) 1.5 C; (d) 2 C. In summary, electrode SOC differences and current density differences combine to cause temperature inhomogeneity, due to the distribution of large-size cell lugs. Electrode SOC and current density inhomogeneity characteristics can be introduced from the temperature inhomogeneity distribution, while electrode SOC and current density differences indicate the physical mechanism of the temperature inhomogeneity.

4.4.1. Battery Connection Optimization with Increased Contact Area of Lug Connection

In the original model design, the arrows in Figure 11a show the cross-section of the current entering and exiting the pole lugs. The maximum temperature difference at the 2 C charge cut-off is 8.8 °C. With all other parameters held constant, the maximum temperature difference at the charging cut-off drops to 4.7 °C. As shown in Figure 11b, by increasing the number of cross sections of the current entry in the model (that is, increasing the cross-sectional area), the maximum temperature drops from 51.2 °C to 46.9 °C, indicating that this method is effective in reducing the temperature inhomogeneity and also has a reducing effect on the amount of heat accumulation of the battery's self-generated heat.

4.4.2. Thermal Management Strategy for Adjusting Heat Transfer Coefficients at Different Positions

Without changing the structural design of the original cell model, the heat dissipation at the two lug side positions of the large-size cell is optimized according to the temperature distribution of 25 °C 2 C operating conditions, and the convective heat transfer coefficients on the positive and negative lug sides are increased from 0.4 W-m-2-K-1 and 0.2 W-m-2-K-1 to 0.5 W-m-2-K-1 and 0.4 W-m-2-K-1, respectively. When comparing Figure 12a,b, it can be seen that the temperature on the lug side of the cell is significantly reduced from about 50 °C to about 40 °C. Additionally, the maximum temperature difference is reduced by more than 1 °C.



Figure 11. Scheme of increasing the pole-ear contact area. (**a**,**c**) Original scheme and its temperature distribution; (**b**,**d**) increasing the lug contact area and its temperature distribution.



Figure 12. Effect of improving the heat exchange on the pole-ear side on the three-dimensional temperature distribution. (a) Temperature distribution of the original scheme; (b) Temperature distribution of the added pole-ear side heat transfer; (c) graphical representation of the improved heat transfer.

Therefore, according to the temperature distribution of the large-size battery at different temperature conditions, for instance, at room temperature, high temperature, and low temperature, cooling can be designed at different locations with different heat transfer coefficients. For example, the value of the convective heat transfer coefficient can be increased by increasing the flow rate of the coolant on the lug side and replacing it with a phase change material with better cooling performance at 2 C charging rate at 25 °C to reduce the uneven distribution of temperature and temperature rise in large-size cells. At low temperatures, the middle part of the large cell surface can be heated separately to improve the charging performance at low temperatures, as the heat generation at the bumps has an elevated effect on the temperature on the lug side.

5. Conclusions

In this paper, the homogeneity of large-size blade lithium-ion batteries is investigated. Firstly, various performance tests were carried out on a large-size blade battery to grasp the capacity of large-size batteries, and the optimal interval of battery performance was obtained. Additionally, based on the inhomogeneity temperature distribution data under different working conditions, the inhomogeneity distribution of large-size batteries was initially grasped. On this basis, a three-dimensional thermoelectric coupling model of a large-size cell was developed, and its accuracy was verified under various operating conditions. The results show that the model can accurately describe the electrochemical and thermal properties of large-size cells under different operating conditions and accurately simulate their inhomogeneity distribution. Finally, based on the model at working condition of 2C charging rate at 25 °C, a thermal management optimization strategy of increasing the contact area of the solder lug and adjusting the heat exchange at different positions is proposed. The results of this optimization strategy reveal that the above solutions is not only effective in alleviating the inhomogeneity distribution of large-size cells, but also in optimizing the temperature rise of the cells, which has an important impact on the charging and discharging performance of the cells under different operating conditions.

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