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Article

Interaction of Wu's Slip Features in Bioconvection of Eyring Powell Nanoparticles with Activation Energy

Anas M. Alwatban ¹, Sami Ullah Khan ², Hassan Waqas ³  and Iskander Tlili ^{4,5,*}

¹ Mechanical Engineering Department, College of Engineering, Qassim University, Qassim 52571, Saudi Arabia; a.alwatban@qu.edu.sa

² Department of Mathematics, COMSATS University Islamabad, Sahiwal 57000, Pakistan; sk_iiu@yahoo.com

³ Department of Mathematics, Government College University, Faisalabad 38000, Pakistan; syedhasanwaqas@hotmail.com

⁴ Department for Management of Science and Technology Development, Ton Duc Thang University, Ho Chi Minh City 758307, Vietnam

⁵ Faculty of Applied Sciences, Ton Duc Thang University, Ho Chi Minh City 758307, Vietnam

* Correspondence: iskander.tlili@tdtu.edu.vn

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Abstract: The current continuation aim is to explore the rheological consequences of Eyring Powell nanofluid over a moving surface in the presence of activation energy and thermal radiation. The bioconvection of magnetized nanoparticles is executed with the evaluation of motile microorganism. The most interesting Wu's slip effects are also assumed near the surface. The evaluation of nanoparticles for current flow problems has been examined by using Buongiorno's model. The governing equations for the assumed flow problem are constituted under the boundary layer assumptions. After converting these equations in dimensionless form, the famous shooting technique is executed. A detailed physical significance is searched out in the presence of slip features. The variation of physical quantities, namely velocity, nanoparticles temperature, nano particles concentration, motile microorganism density, skin friction coefficient, local Nusselt number and motile organism density number are observed with detailed physical aspects for various flow controlling parameters.

Keywords: eyring powell nanofluid; activation energy; motile microorganisms; Wu's slip; shooting technique

1. Introduction

In recent years, nanofluids have proven to be a more convent heat transfer media in contrast to the heat transfer liquids. Due to improved thermo-physical features, the nanoparticles are considered as a more constructive resource energy conversion and heat exchanger, which is quite necessitated in mechanical and engineering processes. The interaction of nanoparticles in base liquids can effectively enhance the cost-coefficient systems, which are related to the diverse mechanical and manufacturing processes. Beside this, the involvement of nano-materials in the era of medical sciences has been found to be quite progressive in recent decades, as these nano-sized particles can be more efficient in the diagnoses of different diseases like cancer tissue, artificial lungs, laser technology, chemotherapy, damage of defected cells etc. The foremost focus on thermo-physical properties of nanoparticles was put forward by Choi [1]. Later on, many miscellaneous attempts regarding the interaction of nanoparticles was made by numerous researchers in the past few years. Madhu and Kishan [2] examined the flow of non-Newtonian nanofluid in the stagnation point mixed convection flow over a moving surface. They employed the finite element numerical technique to execute the numerical

treatment of the formulated problem. The enrolment of nanoparticles in the heat storage units and fins for heat exchangers has been determined by Sheikholeslami and co-workers [3]. Gupta et al. [4] incarcerated the features of thermal radiation and chemical reaction in the flow of nanofluid over an inclined configuration. Hayat et al. [5] performed a theoretical based analysis for nanoparticles suspension in Walters B fluid considered over a nonlinear moving geometry, which sustains a variable thickness. Another interesting contribution of the involvement of nanoparticles in circular jets was treated by Turkyilmazoglu [6]. In this analysis, some interesting analytical and numerical approach has been carried out via famous relevant techniques. Tlili et al. [7] simulated a forced convective flow of nanofluid over a horizontal cylinder in the presence of a magnetic force. The rotating flow of moving disk with Maxwell nanofluid in addition to heat absorption and generation features was reflected by Ahmed et al. [8]. Khan et al. [9] found that entropy generation effects the presence of hybrid nanoparticles. Ahmed et al. [10] investigated the flow of hybrid nanoparticles with micropolar fluid in stenosis arteries. The unsteady slip flow of Carreau nanofluid over a moving wedge in the presence of nonlinear thermal radiation features was focused by Khan et al. [11].

Despite the increasing significance of non-Newtonian fluids, comprehensive attempts have been made to visualize the rheology of non-Newtonian fluids because of their industrial and bioengineering applications. Many fluids like saliva, blood, synovial fluid, honey, silly putty, syrup and cosmetics, protracted a complex relationship for shear stress and deformation. The wide ranging significance of such nonlinear fluids insist the scientist to advise various non-Newtonian fluid models in the literature. Among these fluid models, Eyring Powell is one which reduces to viscous behavior at both high and low shear rates. The unique feature of this model is it does not attain any empirical relation but it is originated from kinetic theory of liquids. The shear thinning features can be depicted using this model. The bhumab body blood is treated as Eyring Powell fluid. This model was suggested by Powell and Eyring [12] in 1944. Later, various investigators analyzed the rheological features of this model with different flow configurations [13–19].

Exclusive investigations relating to the bioconvection of nanoparticles are explored in the current decade, since it involves a variety of significance in the science of biological systems and biotechnology. The basic source of this interesting phenomenon is the upward movement of microorganisms on average. The primly attention for the bioconvection of nanoparticles relates the combined features of pattern formulation and stratification density, which are generated due to self-propelled microorganisms, buoyancy forces, and nanoparticles. The microorganisms usually encompassed are the gravitaxis, gravitaxis, and oxytaxis organisms. The main benefit for combining the features of gyrotactic microorganisms and nanoparticles improve the mass transportation and stability of nanoparticles [20,21]. The basic thought of bioconvection was originated by Kuznetsov [22,23]. Siddiqua et al. [24] discussed the numerical computation of a problem that deals with the bioconvection of nanoparticles in wavy cone. Alsaedi et al. [25] directed the interaction of magneto-nanoparticles along with gyrotactic microorganisms over permeable surface. The effect of Navier slip in nanofluid flow with gyrotactic microorganisms has been depicted by Khan et al. [26]. The impact of temperature dependent thermal conductivity and viscosity in bioconvection of nanoparticles immersed in saturated porous media was worked out by Xun et al. [27]. Waqas et al. [28] utilized the microorganisms' effects in generalized second grade fluid over a stretched surface numerically. In another attempt, Waqas et al. [29] used convective boundary assumption for the flow of Maxwell nanofluid with motile microorganisms. Dero et al. [30] captured multiple solutions for bioconvection that involve nanoparticles in the presence of wall blowing and slip features [31–36].

In the current analysis, we aim to explore the bioconvection phenomenon in the steady flow of Eyring Powell nanofluid in presence of thermal radiation, activation energy and Wu's slip (second order) slip features. The analysis has been performed over a stretched surface. The highly nonlinear formulated problem is solved numerically with shooting procedure. The insight observations for the declared problem are graphically explained.

2. Mathematical Modeling

Let us develop an unsteady and mixed convection flow of Eyring Powell nanofluid in presence of gyrostatic motile microorganism over a moving porous surface. The thermophoresis and Brownian features are entertained in the energy and concentration equations by using Buongiorno's nanofluid model. The flow has been generated due to the moving wedge, which is subjected to the slip boundary conditions. The magnetic field effects are captured by employing it normal to the surface. Moreover, in the energy equation, the thermal radiation features are utilized by using famous Rosseland approximation theory. Following the cartesian coordinate system, velocity component u is assumed in x -direction and v component is taken in y -direction. Let T , C and N explore the involvement of nanoparticles temperature, nanoparticles concentration and motile microorganisms respectively. The flow equation under such flow assumptions are written as [19,27,28]:

$$\frac{\partial u}{\partial x} + \frac{\partial v}{\partial y} = 0 \quad (1)$$

$$u \frac{\partial u}{\partial x} + v \frac{\partial u}{\partial y} = \left(v + \frac{1}{\rho \beta^* C} \right) \frac{\partial^2 u}{\partial y^2} - \frac{1}{2\rho \beta^* C^3} \left[\left(\frac{\partial u}{\partial y} \right)^2 \frac{\partial^2 u}{\partial y^2} \right] - \left(\frac{\sigma B_0^2}{\rho} + \frac{v\varphi}{k'} \right) u + \frac{1}{\rho_f} \left[\begin{array}{l} (1 - C_\infty) \rho_f \beta^* g (T - T_\infty) - (\rho_p - \rho_f) g (C - C_\infty) \\ -(n - n_\infty) g \gamma^* (\rho_m - \rho_f) \end{array} \right], \quad (2)$$

$$u \frac{\partial T}{\partial x} + v \frac{\partial T}{\partial y} = \alpha_0 \frac{\partial^2 T}{\partial y^2} + \frac{(\rho c)_p}{(\rho c)_f} \left\{ D_B \frac{\partial T}{\partial y} \frac{\partial C}{\partial y} + \frac{D_T}{T_\infty} \left(\frac{\partial T}{\partial y} \right)^2 \right\} - \frac{1}{(\rho c)_f} \frac{\partial q_r}{\partial y}, \quad (3)$$

$$u \frac{\partial C}{\partial x} + v \frac{\partial C}{\partial y} = D_B \frac{\partial^2 C}{\partial y^2} + \frac{D_T}{T_\infty} \frac{\partial^2 T}{\partial y^2} - Kr^2 (C - C_\infty) \left(\frac{T}{T_\infty} \right)^2 \exp\left(\frac{-E_a}{kT}\right), \quad (4)$$

$$u \frac{\partial N}{\partial x} + v \frac{\partial N}{\partial y} + \frac{b_1 W_c}{(C_w - C_\infty)} \left[\frac{\partial}{\partial y} \left(N \frac{\partial C}{\partial y} \right) \right] = D_m \left(\frac{\partial^2 N}{\partial y^2} \right), \quad (5)$$

The boundary constraints flowing to the above flow equations are expressed as

$$u = u_w + U_{slip}, v = 0, -k \frac{\partial T}{\partial y} = h_f (T_f - T), D_B \frac{\partial C}{\partial y} + \frac{D_T}{T_\infty} \frac{\partial T}{\partial y} = 0, N = N_w \text{ at } y = 0, \quad (6)$$

$$u \rightarrow 0, v \rightarrow 0, T \rightarrow T_\infty, C \rightarrow C_\infty, N \rightarrow N_\infty \text{ as } y \rightarrow \infty, \quad (7)$$

We define slip boundary conditions U_{slip} in Equation (6) as follows

$$U_{slip} = \frac{2}{3} \left(\frac{3 - \alpha l^2}{\alpha} - \frac{3}{2} \frac{1 - l^2}{K_n} \right) \beta \frac{\partial u}{\partial y} - \frac{1}{4} \left[l^4 + \frac{2}{K_n^2} (1 - l^2) \right] \beta^2 \frac{\partial^2 u}{\partial y^2}, \quad (8)$$

$$U_{slip} = A \frac{\partial u}{\partial y} + B \frac{\partial^2 u}{\partial y^2}, \quad (9)$$

where K_n , A , B , β and α are respectively expressed the Knudsen number, constants, free path for molecular mean and momentum coefficient. It is remarked that these slip conditions are termed as Wu's slip conditions [37–42].

Using the Rosseland approximation, we have

$$q_r = -\frac{4\sigma^*}{3k^*} \frac{\partial T^4}{\partial y}, \quad (10)$$

where σ^* is the Stefan Boltzmann constant and k^* is the absorption constant. For further analysis, we assume that the temperature gradients within the flow are small. As an implication of this assumption

the term T^4 appearing in (16) may be linearized about the ambient temperature T_∞ using Taylor series to give

$$T^4 = T_\infty^4 + 4T_\infty^3(T - T_\infty) + 6T_\infty^2(T - T_\infty)^2 + \dots \quad (11)$$

Upon neglecting the higher-order terms, we get

$$T^4 \cong 4TT_\infty^3 - 3T_\infty^4. \quad (12)$$

In view of (16) and (18), Equation (13) takes the form

$$u \frac{\partial T}{\partial x} + v \frac{\partial T}{\partial y} = \left(\alpha_0 + \frac{16\sigma_s T_\infty^3}{3k^*(\rho c)_f} \right) \frac{\partial^2 T}{\partial y^2} + \frac{(\rho c)_p}{(\rho c)_f} \left\{ D_B \frac{\partial T}{\partial y} \frac{\partial C}{\partial y} + \frac{D_T}{T_\infty} \left(\frac{\partial T}{\partial y} \right)^2 \right\}, \quad (13)$$

Before compute the solution of governing equations, first we convert these equations in dimensionless forms by introducing following variables

$$\psi = \sqrt{av}xf(\eta), \quad \eta = \sqrt{\frac{a}{v}}y, \quad \theta(\eta) = \frac{T - T_\infty}{T_w - T_\infty}, \quad \phi(\eta) = \frac{C - C_\infty}{C_w - C_\infty}, \quad \chi(\eta) = \frac{N - N_\infty}{N_w - N_\infty} \quad (14)$$

We introduce the stream function $\psi(x, y)$ as follows

$$u = \frac{\partial \psi}{\partial y}, \quad v = -\frac{\partial \psi}{\partial x}, \quad (15)$$

The inserting above quantities, the flow Equations (2)–(5) are converted in following forms

$$(1 + K)f'''' - (f')^2 + ff'' - Haf' - K(f'')^2 \lambda f'''' + \Gamma(\theta - Rb\phi - Rc\chi) = 0, \quad (16)$$

$$(1 + Rd)\theta'' + Prf\theta' + Pr(Nb\theta'\phi' + Nt\theta'^2) = 0, \quad (17)$$

$$\phi'' + \left(\frac{Nt}{Nb} \right) \theta'' + PrLe f \phi' - PrLe \sigma (1 + \delta\theta)^n \exp\left(\frac{-E}{1 + \delta\theta} \right) \phi = 0, \quad (18)$$

$$\chi'' + Lp\chi' - Pb[\chi'' + \phi''(\chi + \Omega)] = 0, \quad (19)$$

The transmuted boundary conditions are

$$\left. \begin{aligned} f(\eta) = 0, \quad f'(\eta) = 1 + \alpha f''(\eta) + \beta f''''(\eta), \quad \theta'(0) = Bi(\theta(\eta) - 1), \\ Nb\theta'(\eta) + Nt\phi'(\eta) = 0, \quad \chi(\zeta) = 1 \text{ at } \eta = 0, \end{aligned} \right\} \quad (20)$$

$$f' \rightarrow 0, \quad \theta \rightarrow 0, \quad \phi \rightarrow 0, \quad \chi \rightarrow 0 \text{ as } \eta \rightarrow \infty, \quad (21)$$

In above equations $K = 1/\mu\beta^*C$ and $\lambda = x^2b^3/2vC^2$ are dimensionless material fluid parameters, $Ha = v\varphi/k'b + \sigma B_0^2/\rho b$ combined porosity and magnetic parameter, $Rc = \frac{\gamma_1(\rho_m - \rho_f)(N_w - N_\infty)}{\rho_f(1 - C_\infty)T_\infty \beta_1}$ bio-convection Rayleigh number, buoyancy ratio parameter $Rb = \frac{(\rho_p - \rho_f)(C_w - C_\infty)}{T_\infty \beta_1 \rho_f(1 - C_\infty)}$, $\Gamma = \left(\frac{(T_\infty)\beta_1 g_1(1 - C_\infty)}{a(\rho c)_f} \right)$ mixed convection parameter, $Pr = \frac{\nu}{\alpha}$ is the Prandtl number, $Bi = \frac{h_f}{k} \sqrt{\frac{v}{a}}$ Biot parameter, Lewis number $Le = \frac{\alpha}{D_B}$, $Nt = \frac{\tau(T_w - T_\infty)(D_T)}{\rho c_f(v)}$ thermophoresis parameter, $Nb = \frac{\tau(C_w - C_\infty)(D_B)}{\rho c_f(v)}$ Brownian motion parameter, $Rd = \frac{4\sigma T_\infty^3}{k_1^*k}$ radiation parameter, $\Omega = \frac{N_\infty}{N_w - N_\infty}$, bioconvection constant, $Pb = \frac{bW_c}{D_{mn}}$ are Peclet number, $Lp = \frac{\nu}{D_{mn}}$ is the bio-convection Lewis number, while $\alpha = A\sqrt{\frac{a}{v}}$, and $\beta = B\left(\frac{a}{v}\right)$ first order and second order slip parameters respectively.

3. Numerical Solution

We employ the shooting numerical technique to compute the approximate solution of differential equations [43–47]. For this purpose the boundary value problem (12)–(15) is first converted into initial values problem as follows:

$$\begin{aligned}
 f &= z_1, \quad \frac{df}{d\zeta} = z_2, \quad \frac{d^2f}{d\zeta^2} = z_3, \quad \frac{d^3f}{d\zeta^3} = z'_3, \quad \theta = z_4, \quad \frac{d\theta}{d\zeta} = z_5, \quad \frac{d^2\theta}{d\zeta^2} = z'_5, \\
 \phi &= z_6, \quad \frac{d\phi}{d\zeta} = z_7, \quad \frac{d^2\phi}{d\zeta^2} = z'_7, \quad \chi = z_8, \quad \frac{d\chi}{d\zeta} = z_9, \quad \frac{d^2\chi}{d\zeta^2} = z'_9 \\
 z'_3 &= \frac{1}{(1+K-Kz_3)} \left(z_2^2 - z_1z_3 + Haz_2 - \Gamma(z_4 - Rbz_6 - Rcz_8) \right), \\
 z'_5 &= -\frac{Pr}{(1+Rd)} \left(z_1z_5 - 2z_2z_4 + Nbz_5z_7 + Ntz_5^2 \right), \\
 z'_7 &= -PrLez_1z_7 + PrLe\delta(1 + z_5\delta)^n \exp\left(\frac{-E}{1+\delta z_5}\right) - \frac{Nt}{Nb}z'_5, \\
 z'_9 &= -Lpz_1z_9 + Pb(z'_9 + dz_7(z_8 + \Omega))
 \end{aligned}
 \tag{22}$$

Similarly, the boundary conditions attained following forms:

$$z_1(0) = 0, \quad z_2(0) - 1 - \alpha z_3(0) - \beta z'_3(0), \quad z_6(0) = Bi(z_5(0) - 1), \tag{23}$$

$$Nbz_5(0) + Ntz_7(0) = 0, \quad z_8(0) = 1$$

$$z_2(\infty) \rightarrow 0, \quad z_4(\infty) \rightarrow 0, \quad z_5(\infty) \rightarrow 0, \quad z_7(\infty) \rightarrow 0, \quad z_9(\infty) \rightarrow 0 \tag{24}$$

The numerical computations are perfumed until excellent accuracy with 10^{-6} is obtained.

4. Validation of Solution

In order to verify the solution, the obtained results are compared in Table 1 with already available simulations [35,36]. Table 1 shows that a favorable agreement is noted with these studies.

Table 1. Comparison of solution for $f''(0)$ with various values of Ha when $K = \lambda = Rb = Rc = \Gamma = \alpha = \beta = 0$.

<i>Ha</i>	Wubshet Ibrahim [35]	Ali et al. [36]	Present Results
0.0	1.0000	1.0000	1.0000
1.0	1.4142	1.41421	1.4142
5.0	2.4495	2.44948	2.4496

5. Analysis of Results

The aim of this section is to analyze the physical consequences of involved parameters. For this purpose, various graphs are prepared. Following the traditional theoretical scientific contributions for similar analysis, each parameter has assigned some values, while some fixed values have been accompanied to other remaining flow parameters like $Ha = 0.5$, $\Gamma = 0.2$, $\alpha = 0.2$, $\beta = 0.2$, $Rb = 0.2$, $Rc = 0.1$, $Pr = 0.7$, $Bi = 0.2$, $Rd = 0.4$, $Nb = 0.2$, $Nt = 0.2$, $E = 0.5$, $Pb = 0.4$ and $Lp = 0.1$.

Figure 1 reports the graphical results for combine parameter Ha and mixed convection parameter Γ on the distribution of velocity f' . The combine parameter Ha involves the joint features of both magnetic force and porous medium. With the interaction of magnetic force, the movement of fluid particles resisted as magnetic field results Lorentz force. Similarly, the involvement of permeability of porous medium also reduces the particles motion and as a result the velocity distribution retarded. However, the existence of non-Newtonian parameter retains opposite effects i.e., the distribution of velocity increases with Γ .

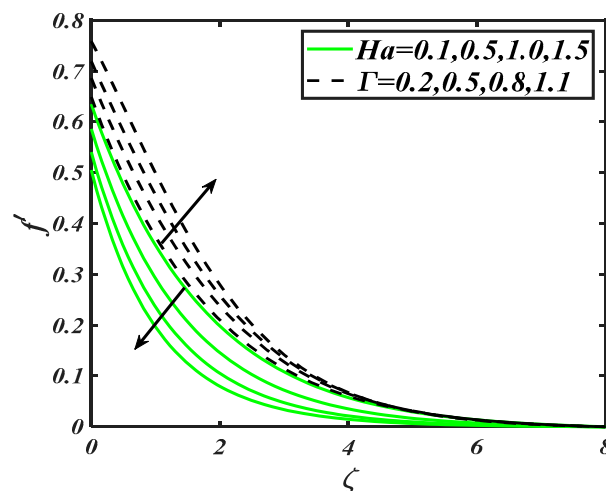


Figure 1. Illustration of Ha and Γ on f' .

The significance of buoyancy ratio parameter Rb and first order slip constant α on f' has been visualized in Figure 2. The observation determined for variation of Rb , shows that the increasing trend in f' is strictly altered effectively. The physical aspect of such an observation may include that bioconvection Rayleigh number involves the buoyancy ratio forces due to which the fluid particles movements decreases. Similarly, the interaction of first order slip also slows down the velocity distribution.

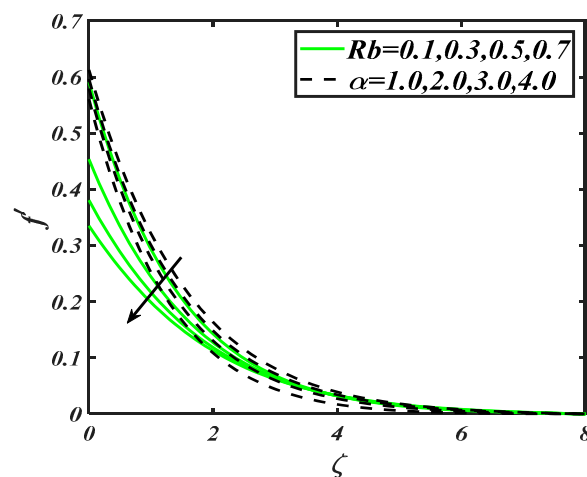


Figure 2. Illustration of Rb and α on f' .

Figure 3 involves the impact of bioconvection Rayleigh number Rc and second order slip factor β on f' . The variation in both parameters depressed the velocity distribution. Physically, evaluation of bioconvection Rayleigh number sustains the buoyancy forces which are of resistive nature. Further, the decreasing trend for velocity in case of β is more dominant as compared to Rc .

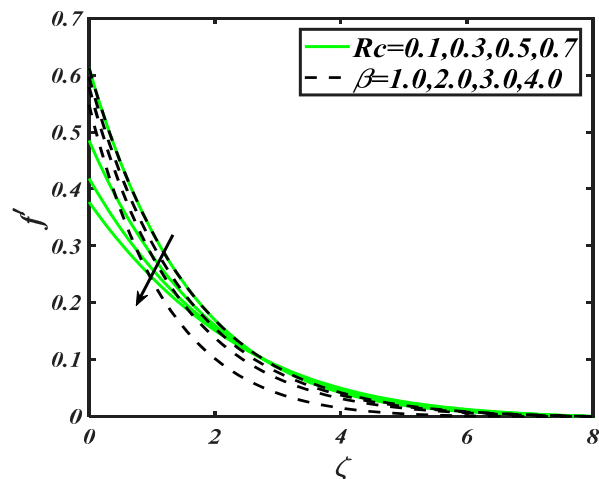


Figure 3. Illustration of Rc and β on f' .

Figure 4 determines the variation in nanoparticles temperature θ against combined parameter Ha and mixed convection parameter Γ . It is observed that the nanoparticles temperature increases linearly with involvement of Ha . Physically, both magnetic force and porous medium permeability enhance the nanoparticles temperature. However, with variation of Γ , the temperature distribution impeded.

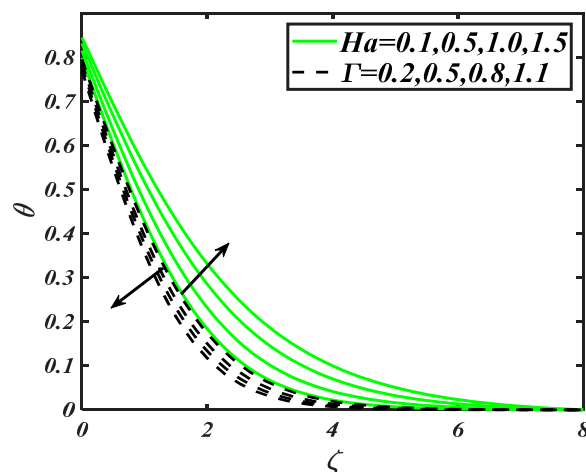


Figure 4. Illustration of Ha and Γ on θ .

Figure 5 perceived the consequences of Prandtl number Pr and radiation constant Rd on θ . A decay in temperature distribution is carried out with variation of Pr . The physical aspects of such a trend may include that Pr is related to the thermal diffusivity inversely due to which θ decreases. On the other hand, increasing values of Rd improve the nanoparticles particles distribution efficiently. Since radiation is considered an active mode of heat transportation, it enables it to raise the temperature of nanoparticles. From Figure 6, it is found that variation in θ reached at maximum level when both thermophoresis constant and thermal Biot number assigned maximum values. Thermophoresis phenomenon involves the migration of fluid particles due to the temperature difference, which yields an enhancement of temperature distribution. Also, thermal Biot number also plays a valuable role in the enhancement of nanoparticles temperature, as it is directly proportional to the heat transfer coefficient due to which θ increases.

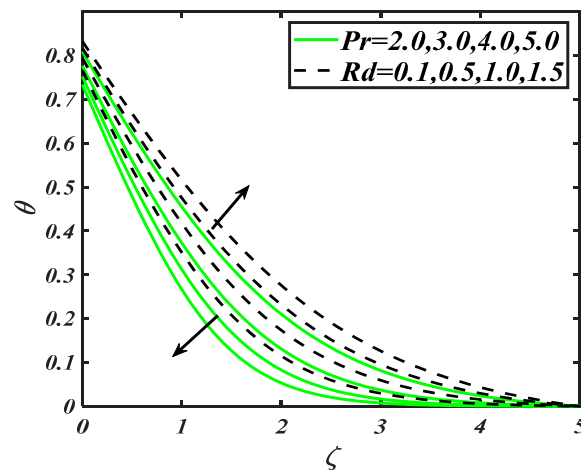


Figure 5. Illustration of Pr and Rd on θ .

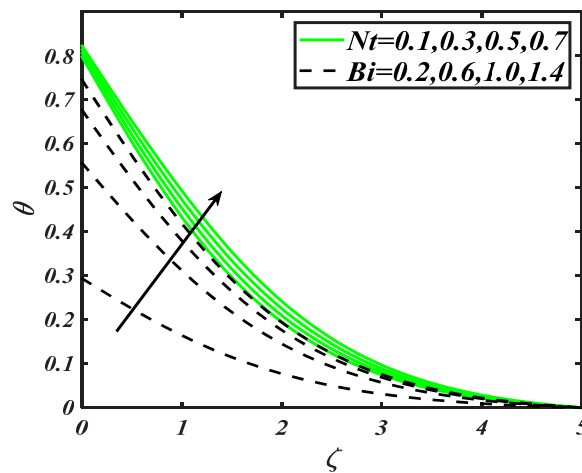


Figure 6. Illustration of Nt and Bi on θ .

To examine the variation in nanoparticles concentration distribution ϕ against flow parameters namely Brownian constant Nb , Lewis number Le , thermophoresis constant Nt , Biot number Bi , activation energy parameter E and Prandtl number Pr , Figures 7–9 are prepared. Figure 7 visualized the graphical features for various prospective of Brownian constant Nb and Lewis number Le . It is securitized that the ϕ decreases with variations of both Nb as Le . Since the Brownian constant reflects the random movement of nanoparticles, which is altered for increasing values Nb . Similarly, the nanoparticles concentration is also impeded for Le . Since Le accomplishes a reverse relationship with mass diffusivity, which causes a reduction in the nanoparticles concentration distribution. From Figure 8, it is noted that a maximum variation in ϕ has been observed, when Nt and Bi assigns maximum values. However, the increasing variation in ϕ is exclusively larger for Nt . The role of activation energy E and Prandtl number Pr constants in distribution of ϕ is notified in Figure 9. With an increase of Pr , the nanoparticles concentration distribution slows down. However, it grows up for activation energy parameter E . Physical features behind this trend concluded the involvement of activation energy provided some extra energy, which enhanced the reaction process.

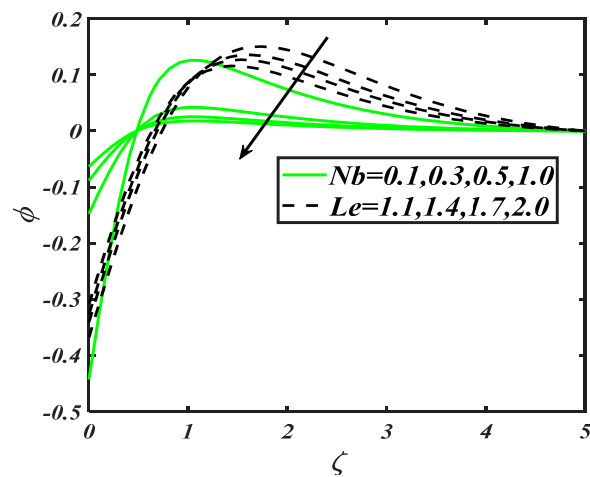


Figure 7. Illustration of Nb and Le on ϕ .

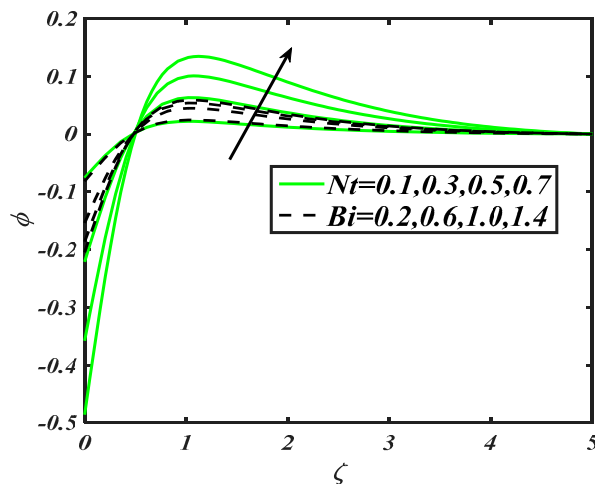


Figure 8. Illustration of Nt and Bi on ϕ .

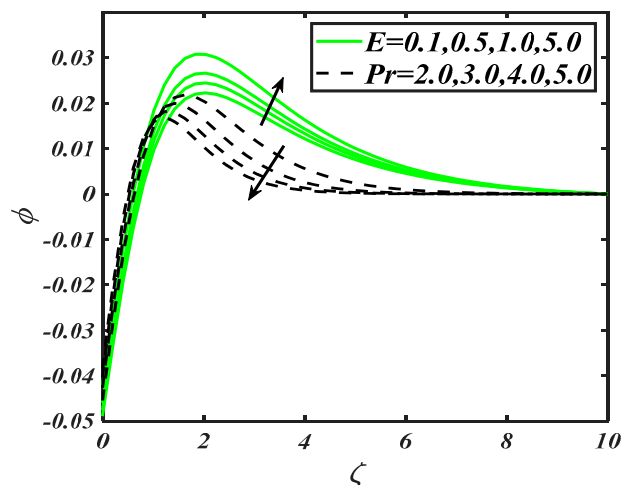


Figure 9. Illustration of E and Pr on ϕ .

The motile microorganisms distribution χ is plotted in Figures 10 and 11 for combined parameter Ha , mixed convection constant Γ , Peclet number Pb and bioconvection Lewis number Lp . The graphical inspected in Figure 10 taken account the effects of Ha and Γ . The motile microorganisms distribution χ raises with increment of Ha . However, in case of Γ , the distribution of motile microorganism increases.

From Figure 11, it is notified that both Peclet number Pb and bioconvection Lewis number Lp decreases the motile microorganism distribution. The dimensionless Peclet number Pb is inversely related to motile diffusivity. Therefore, higher values of Pb accumulate lower motile diffusivity, which turn down the motile microorganisms' distribution.

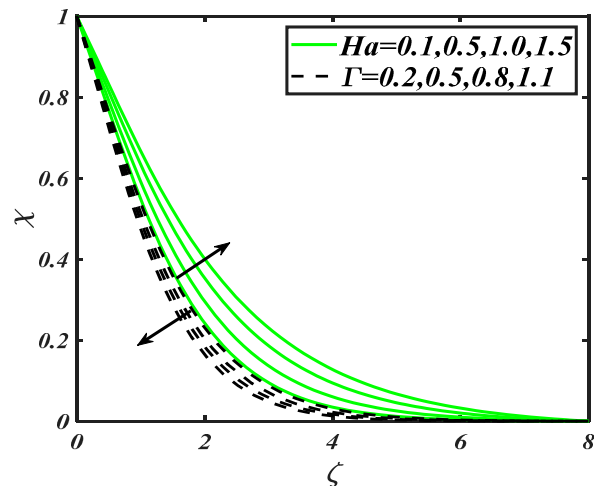


Figure 10. Illustration of Ha and Γ on χ .

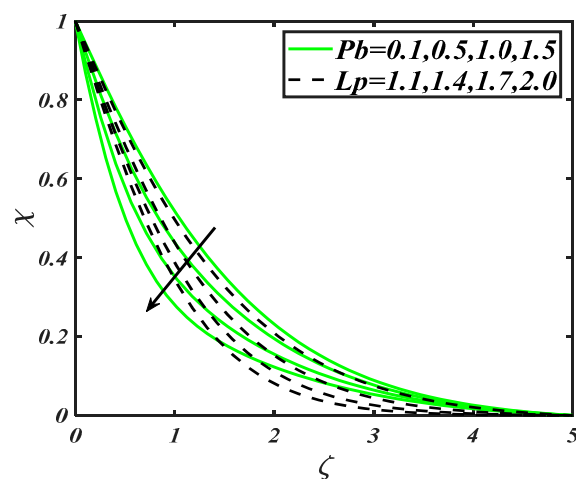


Figure 11. Illustration of Pb and Lp over χ .

Table 2 determines the numerical variation of $-f''(0)$ against values of Ha , K , α , β , Γ , Rb and Rc . The numerical values of $-f''(0)$ increase for Ha , Rb and Rc in contrast to all remaining parameters. The variation in local Nusselt number $-\theta'(0)$ has been captured for varying values of Ha , Rb , Rc , Γ , Pr , Bi , Rd , Nt and Nb in Table 3. Maximum values of local Nusselt number are obtained with variation of Pr while it decreases with all other parameters. From Table 4, we note that the local Sherwood number increases with variation of Pr and Nb . Finally, the variation in local motile density number has been iterated in Table 5. Carefully observing, we see that local motile density number decreases with slip parameters α and β while increasing trend has been noted for Pb and Lp .

Table 2. Variation of skin friction coefficient $-f''(0)$ for $Ha, K, \alpha, \beta, \Gamma, Rb$ and Rc .

Ha	K	α	β	Γ	Rb	Rc	$-f''(0)$
0.1							0.3648
0.4	0.3	1.0	-1.0	0.1	0.1	0.1	0.4027
0.7							0.4335
	0.1						0.4297
	0.4						0.4093
	0.7						0.3924
		2.0					0.2567
		3.0					0.1959
		4.0					1.1589
			-2				0.2489
			-3				0.1872
			-4				0.1532
				0.2			0.3646
				0.5			0.3273
				0.7			0.2941
					0.2		0.3784
					0.3		0.3816
					0.4		0.3876
						0.2	0.3804
						0.3	0.3864
						0.4	0.3927

Table 3. Variation in local Nusselt number $-\theta'(0)$ for $Ha, Rb, Rc, \Gamma, Pr, Bi, Rd, Nt$ and Nb .

Ha	Rb	Rc	Γ	Pr	Bi	Rd	Nt	Nb	$-\theta'(0)$
0.1									0.4022
0.4	0.1	0.1	0.1	0.7	2.0	0.8	0.3	0.2	0.3775
0.7									0.3558
	0.2								0.3935
	0.5								0.3933
	0.8								0.3931
		0.2							0.3923
		0.5							0.3885
		0.8							0.3845
			0.2						0.4016
			0.5						0.4211
			0.8						0.4364
				1.0					0.2865
				3.0					0.4637
				5.0					0.5521
					1.0				0.3306
					1.5				0.3640
					1.5				0.3855
						0.1			0.4848
						0.5			0.4262
						0.8			0.3754
							0.1		0.4104
							0.4		0.3852
							0.7		0.3604
								0.1	0.5521
								0.3	0.3935
								0.4	0.3926

Table 4. Variation in local Sherwood number $-\phi'(0)$ for $\alpha, \beta, \Gamma, Le, Nb, Nt$ and Pr .

α	β	Γ	Le	Nb	Nt	Pr	$-\phi'(0)$
1.0							0.5348
2.0	-1	0.1	0.5	0.2	0.2	0.7	0.5004
3.0							0.4766
	-2.0						0.5307
0.1	-3.0						0.4950
	-4.0						0.4727
		0.2					0.6024
		0.5					0.6316
		0.8					0.6545
			1.0				0.6069
			2.0				0.5995
			3.0				0.5954
				0.2			0.2052
				0.3			0.7703
				0.5			1.2613
					0.4		1.1806
					0.5		0.2952
					0.6		0.1687
						2.0	0.4297
						3.0	0.6956
						4.0	0.8282

Table 5. Variation of motile density number $-\chi'(0)$ for of $Pb, Lp, K, \alpha, \beta, \Gamma, \Omega, Rb$ and Rc .

Pb	Lp	K	α	β	Γ	Ω	Rb	Rc	$-\chi'(0)$
0.2									0.6076
0.6	1.02	0.2	1.0	-1.0	0.1	0.1	0.1	0.1	0.8312
0.8									1.0577
	0.5								0.6115
	1.0								0.7172
	1.5								0.8105
		0.1							0.7759
		0.4							0.8061
		0.7							0.8307
			2.0						0.4911
			3.0						0.4545
			4.0						0.4297
				-1.0					0.4866
				-2.0					0.4488
				-3.0					0.4256
					0.2				0.5657
					0.5				0.5990
					0.8				0.6257
						0.2			0.8150
						0.6			0.8330
						1.0			0.8512
							0.2		0.5521
							0.5		0.5518
							0.8		0.5515
								0.2	0.5502
								0.5	0.5438
								0.8	0.5371

6. Concluding Remarks

We have incorporated the Wu's slip effects, thermal radiation, and activation energy features in bioconvection of Powell Eyring nanofluid over a moving configuration. The formulated problem is solved numerically by incorporating the shooting technique. Some valuable observations from current analysis are summarized as:

- ❖ A devaluate distribution of velocity has been observed for higher values of combine parameter, first order slip parameter and second order slip parameter.
- ❖ The distribution of velocity attains maximum values with mixed convection parameter.
- ❖ The nanoparticles temperature rises with thermophoresis parameter, Biot number and radiation parameter.
- ❖ A decreasing variation in nanoparticles concentration has been figured out for mixed convection parameter and Brownian motion constant.
- ❖ Both Peclet number and the bio-convection Lewis number retarded the motile microorganism distribution.
- ❖ The observations presented here can be simulated to enhance the performances of various thermo-extrusion systems.

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References

1. Choi, S.U.S. *Enhancing Thermal Conductivity of Fluids with Nanoparticles*; No. ANL/MSD/CP-84938; CONF-951135-29; Argonne National Lab.: IL, San Francisco, CA, USA, 1995; Volume 66, pp. 99–105.
2. Madhu, M.; Kishan, N. Finite Element Analysis of heat and mass transfer by MHD mixed convection stagnation-point flow of a non-Newtonian power-law nanofluid towards a stretching surface with radiation. *J. Egypt. Math. Soc.* **2016**, *24*, 458–470. [[CrossRef](#)]
3. Sheikholeslami, M.; Haq, R.; Shafee, A.; Li, Z.; Elaraki, Y.G.; Tlili, I. Heat transfer simulation of heat storage unit with nanoparticles and fins through a heat exchanger. *Int. J. Heat Mass Transf.* **2019**, *135*, 470–478. [[CrossRef](#)]
4. Gupta, S.; Kumar, D.; Singh, J. MHD mixed convective stagnation point flow and heat transfer of an incompressible nanofluid over an inclined stretching sheet with chemical reaction and radiation. *Int. J. Heat Mass Transf.* **2018**, *118*, 378–387. [[CrossRef](#)]
5. Hayat, T.; Qayyum, S.; Alsaedi, A.; Ahmad, B. Magnetohydrodynamic (MHD) nonlinear convective flow of Walters-B nanofluid over a nonlinear stretching sheet with variable thickness. *Int. J. Heat Mass Transf.* **2017**, *110*, 506–514. [[CrossRef](#)]
6. Turkyilmazoglu, M. Free and circular jets cooled by single phase nanofluids. *Eur. J. Mech. B-Fluid* **2019**, *76*, 1–6. [[CrossRef](#)]
7. Tlili, I.; Khan, W.A.; Ramadan, K. MHD Flow of Nanofluid Flow Across Horizontal Circular Cylinder: Steady Forced Convection. *J. Nanofluids* **2019**, *8*, 179–186. [[CrossRef](#)]
8. Ahmed, J.; Khan, M.; Ahmad, L. Stagnation point flow of Maxwell nanofluid over a permeable rotating disk with heat source/sink. *J. Mol. Liq.* **2019**, *287*, 110853. [[CrossRef](#)]
9. Khan, M.I.; Alsaedi, A.; Hayat, T.; Khan, N.B. Modeling and computational analysis of hybrid class nanomaterials subject to entropy generation. *Comput. Meth. Prog. Biol.* **2019**, *179*, 104973. [[CrossRef](#)]
10. Ahmed, A.; Nadeem, S. Effects of magnetohydrodynamics and hybrid nanoparticles on a micropolar fluid with 6-types of stenosis. *Results Phys.* **2017**, *7*, 4130–4139. [[CrossRef](#)]
11. Khan, M.; Azam, M.; Alshomrani, A.S. Unsteady slip flow of Carreaunanofluid over a wedge with nonlinear radiation and new mass flux condition. *Results Phys.* **2017**, *7*, 2261–2270. [[CrossRef](#)]

12. Powell, R.E.; Eyring, H. Mechanisms for the relaxation theory of viscosity. *Nature* **1944**, *154*, 427–428. [[CrossRef](#)]
13. Javed, T.; Ali, N.; Abbas, Z.; Sajid, M. Flow of an Eyring-Powell Non-Newtonian Fluid over a Stretching Sheet. *Chem. Eng. Comm.* **2013**, *200*, 327–336. [[CrossRef](#)]
14. Khan, N.A.; Sultan, F. On the double diffusive convection flow of Eyring-Powell fluid due to cone through a porous medium with Soret and Dufour effects. *AIP Adv.* **2015**, *5*, 057140. [[CrossRef](#)]
15. Khan, S.U.; Ali, N.; Hayat, T. Analytical and Numerical Study of Diffusion of Chemically Reactive Species in Eyring-Powell Fluid over an Oscillatory Stretching Surface. *Bulg. Chem. Commun.* **2017**, *49*, 320–330.
16. Mallick, B.; Misra, J.C. Peristaltic flow of Eyring-Powell nanofluid under the action of an electromagnetic field. *Eng. Sci. Technol. Int. J.* **2019**, *22*, 266–281. [[CrossRef](#)]
17. Akbar, N.S.; Ebaid, A.; Khan, Z.H. Numerical analysis of magnetic field effects on Eyring-Powell fluid flow towards a stretching sheet. *J. Magn. Magn. Mater.* **2015**, *382*, 355–358. [[CrossRef](#)]
18. Malik, M.Y.; Khan, I.; Hussain, A.; Salahuddin, T. Mixed convection flow of MHD Eyring-Powell nanofluid over a stretching sheet: A numerical study. *AIP Adv.* **2015**, *5*, 117118. [[CrossRef](#)]
19. Khan, S.U.; Ali, N.; Abbas, Z. Influence of Heat Generation/Absorption with Convective Heat and Mass Conditions in Unsteady Flow of Eyring Powell Nanofluid over Porous Oscillatory Stretching Surface. *J. Nanofluids* **2016**, *5*, 351–362. [[CrossRef](#)]
20. Wang, S.; Ardekani, A.M. Unsteady swimming of small organisms. *J. Fluid Mech.* **2012**, *702*, 286–297. [[CrossRef](#)]
21. Elfring, G.J.; Goyal, G. The effect of gait on swimming in viscoelastic fluids. *J. Non-Newton Fluid* **2016**, *234*, 8–14. [[CrossRef](#)]
22. Kuznetsov, A.V. The onset of nanofluidbioconvection in a suspension containing both nanoparticles and gyrotactic microorganisms. *Int. Commun. Heat Mass Transf.* **2010**, *37*, 1421–1425. [[CrossRef](#)]
23. Kuznetsov, A.V. Nanofluidbioconvection in water-based suspensions containing nanoparticles and oxytactic microorganisms: Oscillatory instability. *Nanoscale Res. Lett.* **2011**, *6*, 100. [[CrossRef](#)] [[PubMed](#)]
24. Siddiqua, S.; Begum, N.; Saleem, S.; Hossain, M.A.; Gorla, R.S.R. Numerical solutions of nanofluid bio-convection due to gyrotactic microorganisms along a vertical wavy cone. *Int. J. Heat Mass Transf.* **2016**, *101*, 608–613. [[CrossRef](#)]
25. Khan, A.; Farooq, M.; Gull, N.; Hayat, T. Magneto hydrodynamic (MHD) stratified bio-convective flow of nanofluid due to gyrotactic microorganisms. *Adv. Powder Technol.* **2017**, *28*, 288–298.
26. Khan, W.A.; Uddin, M.J.; Ismail, A.I.M. Free convection of non-Newtonian nanofluids in porous media with gyrotactic microorganisms. *Trans. Porous Media* **2013**, *97*, 241–252. [[CrossRef](#)]
27. Xun, S.; Zhao, J.; Zheng, L.; Zhang, X. Bio convection in rotating system immersed in nanofluid with temperature dependent viscosity and thermal conductivity. *Int. J. Heat Mass Transf.* **2017**, *111*, 1001–1006. [[CrossRef](#)]
28. Waqas, H.; Khan, S.U.; Hassan, M.; Bhatti, M.M.; Imran, M. Analysis on the bio convection flow of modified second-grade nanofluid containing gyrotactic microorganisms and nanoparticles. *J. Mol. Liq.* **2019**, *291*, 111231. [[CrossRef](#)]
29. Waqas, H.; Khan, S.U.; Shehzad, S.A.; Imran, M. Radiative flow of Maxwell nanofluid containing gyrotactic microorganisms and energy activation with convective Nield conditions. *Heat Transf.-Asian Res.* **2019**, *48*, 1663–1687. [[CrossRef](#)]
30. Dero, S.; Uddin Md, J.; Rohni, A.M.R. Stefan Blowing and Slip Effects on Unsteady Nanofluid Transport Past a Shrinking Sheet: Multiple Solutions. *Heat Transf. Asian Res.* **2019**. [[CrossRef](#)]
31. Tlili, I. Effects MHD and Heat Generation on Mixed Convection Flow of Jeffrey Fluid in Microgravity Environment over an Inclined Stretching Sheet. *Symmetry* **2019**, *11*, 438. [[CrossRef](#)]
32. Goodarzi, M.; Tlili, I.; Tian, Z.; Safaei, M. Efficiency assessment of using graphene nanoplatelets-silver/water nanofluids in microchannel heat sinks with different cross-sections for electronics cooling. *Int. J. Numer. Methods Heat Fluid Flow* **2019**. [[CrossRef](#)]
33. Afridi, M.I.; Tlili, I.; Goodarzi, M.; Osman, M.; Khan, N.A. Irreversibility Analysis of Hybrid Nanofluid Flow over a Thin Needle with Effects of Energy Dissipation. *Symmetry* **2019**, *11*, 663. [[CrossRef](#)]
34. Tlili, I. Tawfeeq Abdullah Alkanhal, Nanotechnology for water purification: Electrospun nanofibrous membrane in water and wastewater treatment. *J. Water Reuse Desalin.* **2019**, *24*. [[CrossRef](#)]

35. Tlili, Iskander. Thermodynamic Study on Optimal Solar Stirling Engine Cycle Taking Into Account the Irreversibilities Effects. *Energy Procedia*. **2012**, 584–591. [[CrossRef](#)]
36. Tlili, I.; Khan, W.A.; Ramadan, K. Entropy Generation Due to MHD Stagnation Point Flow of a Nanofluid on a Stretching Surface in the Presence of Radiation. *J. Nanofluids* **2018**, *7*, 879–890. [[CrossRef](#)]
37. Wu, L. A slip model for rarefied gas flows at arbitrary Knudsen number. *Appl. Phys. Lett.* **2008**, *93*, 253103. [[CrossRef](#)]
38. Nandeppanavar, M.M.; Vajravelu, K.; Abel, M.S.; Siddalingappa, M. Second order slip flow and heat transfer over a stretching sheet with non-linear Navier boundary condition. *Int. J. Therm. Sci.* **2012**, *58*, 143–150. [[CrossRef](#)]
39. Rosca, A.V.; Pop, I. Flow and heat transfer over a vertical permeable stretching/shrinking sheet with a second order slip. *Int. J. Heat Mass Transf.* **2013**, *60*, 355–364. [[CrossRef](#)]
40. Khan, S.U.; Waqas, H.; Shehzad, S.A.; Imran, M. Theoretical analysis for tangent hyperbolic nanoparticles with combined electrical MHD, activation energy and Wu's slip features: A mathematical model. *Phys. Scr.* **2019**. [[CrossRef](#)]
41. Ibrahim, W. Magnetohydrodynamics (MHD) flow of a tangent hyperbolic fluid with nanoparticles past a stretching sheet with second order slip and convective boundary condition. *Results Phys.* **2017**, *7*, 3723–3731. [[CrossRef](#)]
42. Ali, F.M.; Nazar, R.; Arifin, N.M.; Pop, I. MHD boundary layer flow and heat transfer over a stretching sheet with induced magnetic field. *Heat Mass Transf.* **2011**, *47*, 155–162. [[CrossRef](#)]
43. Tlili, I.; Khan, W.A.; Khan, I. Multiple slips effects on MHD SA-Al₂O₃ and SA-Cu non-Newtonian nanofluids flow over a stretching cylinder in porous medium with radiation and chemical reaction. *Results Phys.* **2018**, *8*, 213–222. [[CrossRef](#)]
44. Tlili, I.; Hamadneh, N.N.; Khan, W.A. Thermodynamic Analysis of MHD Heat and Mass Transfer of Nanofluids Past a Static Wedge with Navier Slip and Convective Boundary Conditions. *Arab. J. Sci. Eng.* **2018**, 1–13. [[CrossRef](#)]
45. Tlili, I.; Hamadneh, N.N.; Khan, W.A.; Atawneh, S. Thermodynamic analysis of MHD Couette—Poiseuille flow of water—Based nanofluids in a rotating channel with radiation and Hall effects. *J. Therm. Anal. Calorim.* **2018**, *132*, 1899–1912. [[CrossRef](#)]
46. Khan, W.A.; Rashad, A.M.; Abdou, M.M.M.; Tlili, I. Natural bioconvection flow of a nanofluid containing gyrotactic microorganisms about a truncated cone. *Eur. J. Mech.-B/Fluids* **2019**, *75*, 133–142. [[CrossRef](#)]
47. Khalid, A.; Khan, I.; Khan, A.; Shafie, S.; Tlili, I. Case study of MHD blood flow in a porous medium with CNTs and thermal analysis. *Case Stud. Therm. Eng.* **2018**, *12*, 374–380. [[CrossRef](#)]



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