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Adaptive proportional resonant (PR) controllers, whose resonant frequencies are obtained by the phase-locked loop (PLL), are employed in grid connected voltage source converters (VSCs) to improve the control performance in the case of grid frequency variations. The resonant frequencies can be estimated by either synchronous reference frame PLL (SRF-PLL) or dual second order generalized integrator frequency locked loop (DSOGI-FLL), and there are three different implementations of the PR controllers based on two integrators. Hence, in this paper, system stabilities of the VSC with different implementations of PR controllers and different PLLs under weak grid conditions are analyzed and compared by applying the impedance-based method. First, the s -domain admittance matrixes of the VSC are derived using the harmonic linearization method. Then, the admittance matrixes are compared with each other, and the influences of their differences on system stability are revealed. It is demonstrated that if DSOGI-FLL is used, stabilities of the VSC with different implementations of the PR controllers are similar. Moreover, the VSC using a DSOGI-FLL is more stable than that using a SRF-PLL. The simulation and experimental results are conducted to verify the correctness of theoretical analysis.

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Article

Stability Analysis of Grid-Connected Converters with Different Implementations of Adaptive PR Controllers under Weak Grid Conditions

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Abstract: Adaptive proportional resonant (PR) controllers, whose resonant frequencies are obtained by the phase-locked loop (PLL), are employed in grid connected voltage source converters (VSCs) to improve the control performance in the case of grid frequency variations. The resonant frequencies can be estimated by either synchronous reference frame PLL (SRF-PLL) or dual second order generalized integrator frequency locked loop (DSOGI-FLL), and there are three different implementations of the PR controllers based on two integrators. Hence, in this paper, system stabilities of the VSC with different implementations of PR controllers and different PLLs under weak grid conditions are analyzed and compared by applying the impedance-based method. First, the $\alpha\beta$ -domain admittance matrixes of the VSC are derived using the harmonic linearization method. Then, the admittance matrixes are compared with each other, and the influences of their differences on system stability are revealed. It is demonstrated that if DSOGI-FLL is used, stabilities of the VSC with different implementations of the PR controllers are similar. Moreover, the VSC using a DSOGI-FLL is more stable than that using a SRF-PLL. The simulation and experimental results are conducted to verify the correctness of theoretical analysis.

Keywords: grid-connected converter; adaptive resonant controller; PLL; impedance analysis

1. Introduction

With the increasing energy consumption worldwide, the use of renewable energy sources like wind and solar energies [1,2] in the grid has been growing increasingly. The voltage source converters (VSCs) have many desirable features, such as full controllability, low current harmonics, and high efficiency, thus they are widely used to deliver the power produced by the renewable energies into the grid [3].

Proportional resonant (PR) controllers could control positive-sequence and corresponding negative-sequence grid current at the same time without any additional negative-sequence current controller, and they allow a relatively low computational cost as they are implemented in the stationary frame [3–5]. Hence, PR controllers working on the stationary reference frame are widely used in grid-connected VSCs. Implementations of the PR controllers based on two integrators are widely employed, since no explicit trigonometric functions are needed [5]. Three typical implementations of the resonance term of the current controllers in the prior studies are shown in Figure 1 [5–7]. For future reference, they are called implementation I, II and III, respectively. If the resonant frequency is constant, these implementations are equivalent to each other.

The grid frequency is practically not a constant but within a certain range [8,9], thus the control performance would be inevitable weakened if the center frequency of the resonance controller is set to

be constant. If the resonant frequency is set to be the fundamental frequency and the grid frequency deviates from it, there would be a phase shift between the grid current and the corresponding grid voltage, though the aimed power factor is unity [6]. To improve control performance in the case of grid frequency variations, the authors of [6,8,10–13] developed frequency adaptive PR controllers, whose resonant frequencies are not constant values, but are updated online according to the frequency estimated by the phase-locked loop (PLL) system. If the adaptive PR controllers are applied, a unity power factor operation can be achieved [6]. The resonant frequency of the adaptive PR controller is time-variant, thus dynamic properties of the adaptive PR controllers with different implementations might be dramatically different. Consequently, the port characteristics of the converter with different implementations of adaptive PR controllers would not be the same.

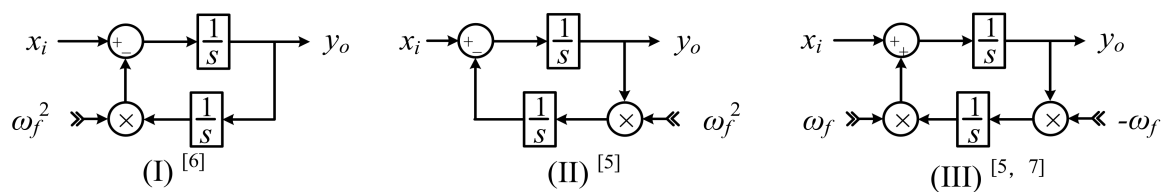


Figure 1. Implementations of the resonant term of frequency adaptive proportional resonant (PR) controllers based on two integrators.

The resonant frequency of the adaptive PR controller can be obtained by grid synchronization techniques based on phase locking approach, e.g., synchronous reference frame-phase locked loop (SRF-PLL) [14] which is widely used for its simplicity and robustness, or frequency locking approach, e.g., dual second-order generalized integrator-frequency locked loop (DSOGI-FLL) [15] which is implemented in the stationary reference frame. As SRF-PLL and DSOGI-FLL use different approaches to obtain the grid frequency, their influences on the port characteristics of the converter with adaptive current controllers might differ be different.

Under weak grid conditions, stability issues introduced by the grid-connected VSCs are of great importance [16–22]. The differences of the port characteristics of the converter introduced by applying different control schemes, including three different implementations of the adaptive PR controllers and the two different PLLs to obtain the resonant frequency, might significantly change the stability of the VSC connected to a weak grid. Hence, it is necessary to compare the robustness of the VSC with different control schemes.

In this paper, the stability issues of the VSC with different implementations of the adaptive PR controllers and different PLLs are studied and analyzed using the impedance-based method [18,23–28]. A suggestion to choose a suitable controllers' implementation and the related PLL for the VSC with adaptive PR controllers is given.

The rest of the paper is organized as follows: in Section 2, the studied system with adaptive current controller is briefly introduced. Section 3 shows how to model the adaptive resonant controller with different implementations of the resonant terms, and the approach to incorporate it into the admittance model of the converter. In Section 4 the effects of the adaptive resonant controllers with different implementations on system stability is analyzed and compared, and the stability of the system using a SRF-PLL for grid synchronization and the system using a DSOGI-FLL for grid synchronization is compared. Section 5 includes experimental verifications of the theoretical analysis. Section 6 concludes this paper.

2. VSC with Adaptive PR Current Controllers

The L-type grid connected converter with grid current regulation working on $\alpha\beta$ reference frame is studied, as depicted in Figure 2. The inductance of the filter is L , and the equivalent series resistance is r . The grid inductance is L_g . The grid voltages, the voltages at the point of common coupling (PCC),

the ac-side converter voltages, and the currents delivered to the grid are u_{sk} , u_{gk} , v_k , and i_k ($k = a, b, c$), respectively. The dc input voltage of the converter is V_{dc} .

Grid current references in the stationary reference frame ($i_{\alpha r}$, $i_{\beta r}$) can be obtained by applying an inverse Park transformation to the active, reactive current references (I_{dr} , I_{qr}). In this paper, unit-power factor is considered, i.e., I_{qr} is set to be zero. Subscripts ‘ α ’ and ‘ β ’ refer to the α -axis and β -axis, respectively, while subscripts ‘ d ’ and ‘ q ’ refer to the d -axis and q -axis, respectively.

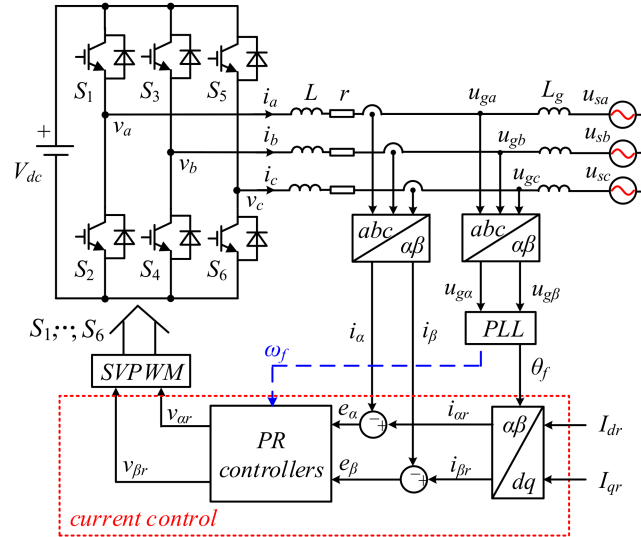


Figure 2. Block diagram of a voltage source converter (VSC) with adaptive PR controllers for grid-connected applications.

2.1. Adaptive PR Controller

The grid current error signals (e_α , e_β) are sent to the PR controllers to generate the reference of the ac-side converter voltages ($v_{\alpha r}$, $v_{\beta r}$). The PR current controllers are defined as follows:

$$H(s) = k_p + k_r \frac{s}{s^2 + \omega_f^2} = k_p + k_r R(s) \tag{1}$$

where ω_f is the resonant frequency, different from the conventional PR controller, the adaptive PR controller uses the frequency estimated by the PLL as the resonant frequency instead of the constant fundamental frequency ω_1 ; k_p , k_r are the proportional- and resonant-gain of the controller, respectively. The controller gains can be tuned according to: $k_p = \alpha_c L$, $k_i = \alpha_c r$ [18] where α_c is the current control loop bandwidth. $R(s)$ is the resonant term of the controller. $R(s)$ has infinite gain at the resonant frequency and thus it is capable for the grid current to track its reference without steady-state error.

The detailed block diagrams of different implementations of the PR controller based on two integrators are depicted in Figure 3, where $x_{i\alpha}$ and $y_{o\alpha}$ are the input and output of $R(s)$, respectively. The α -axis and β -axis are decoupled, thus only the PR controller implemented in the α -axis is given for simplicity. In time domain, the relationships between $x_{i\alpha}$ and $y_{o\alpha}$ can be derived from Figure 3, as follows:

$$\begin{cases} x_{i\alpha}(t) = \frac{dy_{o\alpha}(t)}{dt} + \omega_f^2(t) \int_0^t y_{o\alpha}(t) dt \cdots \text{I} \\ x_{i\alpha}(t) = \frac{dy_{o\alpha}(t)}{dt} + \int_0^t \omega_f^2(t) y_{o\alpha}(t) dt \cdots \text{II} \\ x_{i\alpha}(t) = \frac{dy_{o\alpha}(t)}{dt} + \omega_f(t) \int_0^t \omega_f(t) y_{o\alpha}(t) dt \cdots \text{III} \end{cases} \tag{2}$$

where Equations (2-I), (2-II) and (2-III) are obtained with implementation I, II, and III of the PR controllers, respectively. If ω_f is constant, the three equations are equivalent to each other. However, for the adaptive PR controllers, ω_f is time varying, the equations are no longer the same. Thus the

dynamic properties of the adaptive PR controllers with different implementations might be different, and the transfer function shown in Equation (1) is not enough to describe the dynamic properties of adaptive PR controllers.

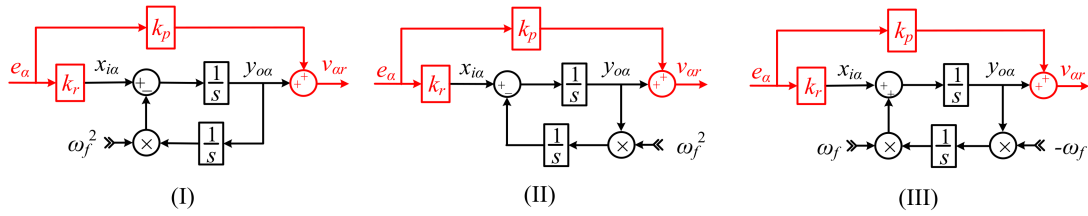


Figure 3. Detailed block diagrams of adaptive PR controllers with different implementations of $R(s)$.

The adaptive PR controllers’ outputs are sent to the SVPWM to generate the control signals.

2.2. Grid Synchronization Methods

Two typical grid synchronization methods, i.e., SRF-PLL and DSOGI-FLL as shown in Figure 4 are used to provide the phase angle to calculate the reference currents and the controller’s resonant frequency. The block diagram of a typical SRF-PLL is shown in Figure 4a, where $G_{PLL}(s)$ is the PLL controller. Figure 4b shows the block diagram of DSOGI-FLL with frequency locked loop (FLL) gain normalization, where k is the gain of the SOGI, and γ is used to regulate the settling time of the FLL [22]. The DSOGI-FLL consists of four parts: (1) SOGI; (2) positive/negative-sequence calculation block (PNSC) to obtain the positive-sequence voltages ($u_{\alpha+}$, $u_{\beta+}$) and the negative-sequence voltage ($u_{\alpha-}$, $u_{\beta-}$); (3) FLL with FLL gain normalization to estimate the grid frequency; (4) the “atan2” uses the positive-sequence voltages for the phase angle calculation.

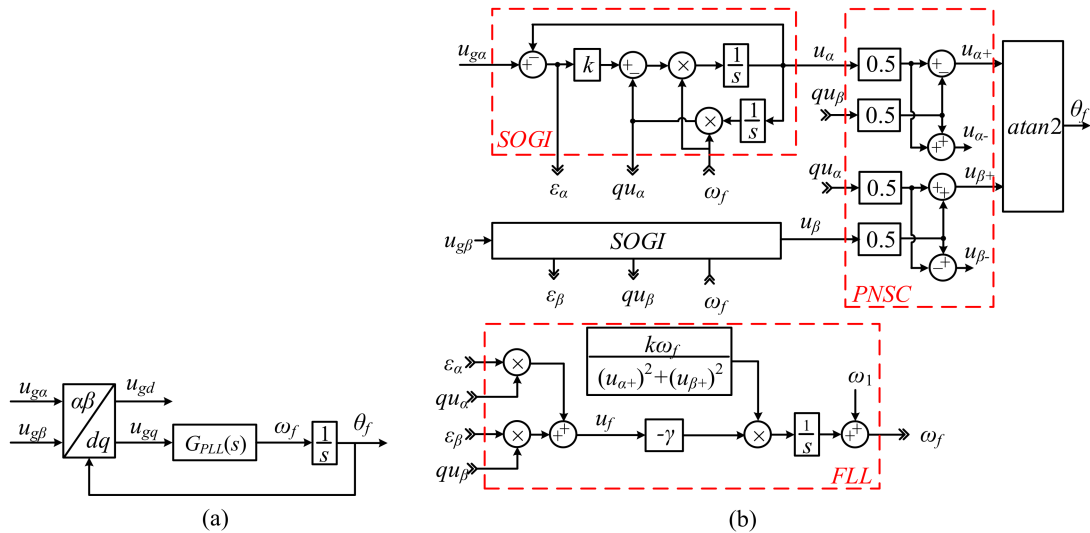


Figure 4. Block diagram of (a) a basic synchronous reference frame PLL (SRF-PLL); (b) a dual second order generalized integrator frequency locked loop (DSOGI-FLL) based grid synchronization with frequency locked loop (FLL) gain normalization.

With different implementations of the adaptive PR controllers and different PLLs, the stability of the grid-connected converter might be quite different under weak grid conditions.

3. Impedance Modeling

Impedance-based methods were a popular choice in prior studies to analyse the stability of grid-connected converters. Impedance modeling of the converter can be performed in either the dq -domain [25] or in the $\alpha\beta$ -domain [26,27]. Considering that the grid current regulation working on $\alpha\beta$ reference frame is used, impedance-based method in the $\alpha\beta$ -domain is applied in this paper. The $\alpha\beta$ -domain converter impedances are derived using harmonic linearization method.

A positive-sequence perturbation at an arbitrary frequency ω_p is injected to the grid to obtain the reflected admittances at the ac terminals. Thus, in three phase variables, there would have positive-sequence components at frequency ω_p and negative-sequence components at frequency $\omega_p - 2\omega_1$ [28]. In this paper, it is defined that $s = j\omega$, where $\omega = \omega_p - \omega_1$.

3.1. Model of the Adaptive PR Controllers

The outputs of the adaptive PR controllers are affected by the error signal and the estimated frequency, as shown in Figure 3, thus the small-signal part of controllers' outputs can be described by Equation (3):

$$\begin{cases} v_{\alpha r}^p(\omega + \omega_1) = H(s + j\omega_1)e_{\alpha}^p(\omega + \omega_1) + H_p^x(s + j\omega_1)\omega_f(\omega) \\ v_{\alpha r}^n(\omega - \omega_1) = H(s - j\omega_1)e_{\alpha}^n(\omega - \omega_1) + H_n^x(s - j\omega_1)\omega_f(\omega) \end{cases} \quad (2)$$

where $H(s)$ is defined in Equation (1); $H_p^x(s)$ and $H_n^x(s)$ are used to describe the influence of ω_f on the positive-, negative-sequence components of the current controllers' outputs, respectively. The superscripts of the variables 'p' and 'n' are used to indicate the positive-, negative-sequence of the variables, respectively. The superscripts of the transfer functions 'x' in $H_p^x(s)$ and $H_n^x(s)$ is used to distinguish the implementations of adaptive controllers: $x = i, ii, \text{ and } iii$ indicates that implementation I, II and III is used, respectively.

To obtain the detailed expressions of $H_p^x(s)$ and $H_n^x(s)$, e_{α} is assumed to be zero, thus from Figure 3, it can be easily obtained that $x_{i\alpha} = 0$, and $y_{o\alpha} = v_{\alpha r}$. If the implementation I of the adaptive PR controller is used, it can be derived from (2-I) by using convolution that:

$$\begin{cases} (s + j\omega_1)v_{\alpha r}^p(\omega + \omega_1) + \left(\omega_1^2 \frac{v_{\alpha r}^p(\omega + \omega_1)}{s + j\omega_1} + 2\omega_1\omega_f(\omega) \frac{V_{\alpha}[\omega_1]}{j\omega_1} \right) = 0 \\ (s - j\omega_1)v_{\alpha r}^n(\omega - \omega_1) + \left(\omega_1^2 \frac{v_{\alpha r}^n(\omega - \omega_1)}{s - j\omega_1} + 2\omega_1\omega_f(\omega) \frac{V_{\alpha}[-\omega_1]}{-j\omega_1} \right) = 0 \end{cases} \quad (3)$$

where $V_{\alpha}[\pm\omega_1] = (V_m + rI_{dr} \pm jL\omega_1 I_{dr})/2$ are the Fourier coefficients of $v_{\alpha r}$ at the positive/negative fundamental frequency, V_m is the magnitude of PCC voltage. For simplicity, set $V_p = 2V_{\alpha}[\omega_1]$, and $V_n = 2V_{\alpha}[-\omega_1]$.

The detailed expressions of $H_p^x(s)$ and $H_n^x(s)$ for $x = i$ can be easily obtained from Equations (3) and (4), as shown in Table 1. Following the same procedure, $H_p^x(s)$ and $H_n^x(s)$ for $x = ii, \text{ and } iii$ can also be derived as shown in Table 1.

Table 1. Detailed expressions of $H_p^x(s)$ and $H_n^x(s)$ in Equation (3) with different implementations of the adaptive PR current controllers.

X	i	ii	iii
$H_p^x(s + j\omega_1)$	$\frac{jV_p(s + j\omega_1)}{s(s + j2\omega_1)}$	$-\frac{V_p\omega_1}{s(s + j2\omega_1)}$	$j\frac{V_p}{2s}$
$H_n^x(s - j\omega_1)$	$-\frac{jV_n(s - j\omega_1)}{s(s - j2\omega_1)}$	$-\frac{V_n\omega_1}{s(s - j2\omega_1)}$	$-j\frac{V_n}{2s}$

3.2. Model of the PLL System

The frequency domain forms of ω_f and θ_f are shown in Equation (5). The superscripts 'y' is used to distinguish the PLL that is applied: $y = pll$ means that a SRF-PLL is used; $y = fl$ means that a DSOGI-FLL is used. The expressions of the transfer functions are shown in Table 2 where

$T_{pll}(s) = G_{PLL}(s)/(s + V_m G_{PLL}(s))$ and $D(s) = s^2 + k\omega_1 s + \omega_1^2$. The derivations are shown in the Appendix A.

$$\begin{cases} \omega_f(\omega) = H_{wp}^y(s)u_{g\alpha}^p(\omega + \omega_1) + H_{wn}^y(s)u_{g\alpha}^n(\omega - \omega_1) \\ \theta_f(\omega) = H_{sp}^y(s)u_{g\alpha}^p(\omega + \omega_1) + H_{sn}^y(s)u_{g\alpha}^n(\omega - \omega_1) \end{cases} \quad (4)$$

Table 2. Detailed expressions of the transfer functions in Equation (5) with different phase-locked loops (PLLs).

$H_{sp}^{pll}(s) = -jT_{pll}(s)$	$H_{wp}^{pll}(s) = -jsT_{pll}(s)$
$H_{sn}^{pll}(s) = jT_{pll}(s)$	$H_{wn}^{pll}(s) = jsT_{pll}(s)$
$H_{sp}^{fll}(s) = -\frac{j k\omega_1(s+j2\omega_1)}{2V_m D(s+j\omega_1)} \left(1 + \frac{\lambda \left(\frac{s-j2\omega_1}{D(s-j\omega_1)} + \frac{s+j2\omega_1}{D(s+j\omega_1)} \right)}{1 + \lambda \frac{k\omega_1}{s} \left(\frac{s-j\omega_1}{D(s-j\omega_1)} + \frac{s+j\omega_1}{D(s+j\omega_1)} \right)} \right)$	$H_{wp}^{fll}(s) = \frac{-j\lambda \frac{k\omega_1}{V_m} \frac{s+j2\omega_1}{D(s+j\omega_1)}}{1 + \lambda \frac{k\omega_1}{s} \left(\frac{s-j\omega_1}{D(s-j\omega_1)} + \frac{s+j\omega_1}{D(s+j\omega_1)} \right)}$
$H_{sn}^{fll}(s) = j \frac{k\omega_1(s-j2\omega_1)}{2V_m D(s-j\omega_1)} \left(1 + \frac{\lambda \left(\frac{s-j2\omega_1}{D(s-j\omega_1)} + \frac{s+j2\omega_1}{D(s+j\omega_1)} \right)}{1 + \lambda \frac{k\omega_1}{s} \left(\frac{s-j\omega_1}{D(s-j\omega_1)} + \frac{s+j\omega_1}{D(s+j\omega_1)} \right)} \right)$	$H_{wn}^{fll}(s) = \frac{j\lambda \frac{k\omega_1}{V_m} \frac{s-j2\omega_1}{D(s-j\omega_1)}}{1 + \lambda \frac{k\omega_1}{s} \left(\frac{s-j\omega_1}{D(s-j\omega_1)} + \frac{s+j\omega_1}{D(s+j\omega_1)} \right)}$

3.3. Model of the Current References

The $\alpha\beta$ -domain current references in the frequency domain is shown in Equation (6) [22]:

$$\begin{cases} i_{ar}^p(\omega + \omega_1) = j0.5I_{dr}\theta_f(\omega) \\ i_{ar}^n(\omega - \omega_1) = -j0.5I_{dr}\theta_f(\omega) \end{cases} \quad (5)$$

3.4. Model of the Grid Current Loop

Based on the adaptive PR controller (3), the PLL system (5) and the current references (6), the detailed small-signal model of the grid current with adaptive PR controllers can be obtained, as depicted in Figure 5, where $G_d(s)$ is used to depicts the gain and delays (digital computation delay and the PWM delay) [29] of the converter, as follows:

$$G_d(s) = e^{-1.5sT_s} \quad (6)$$

where T_s is the sampling period.

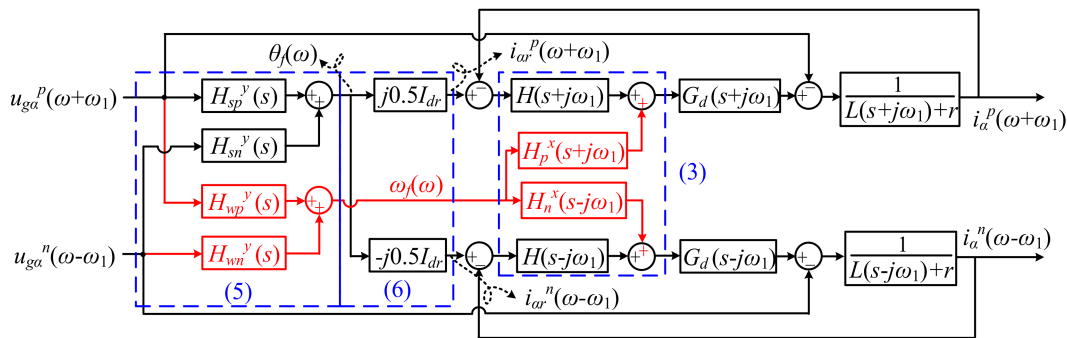


Figure 5. Small signal model of grid current control loop influenced by the PLL system.

Because of the adaptive PR controller, extra branches are introduced in the model of grid current loop, as shown in the red part of Figure 5. Moreover, $H_p^x(s)$, $H_n^x(s)$ (as shown in Table 1) are different for different implementations of adaptive PR controllers, and $H_{wp}^y(s)$, $H_{wn}^y(s)$ (as shown in Table 2) are

not the same for different PLLs. Hence, the port characteristics of the grid-connected converters with different implementations of the adaptive PR controllers and different PLLs are different.

The small-signal part of grid current can be obtained from Figure 5, as follows:

$$\begin{cases} i_{\alpha}^p(\omega + \omega_1) = T_i(s + j\omega_1)i_{\alpha r}^p(\omega + \omega_1) - Y_i(s + j\omega_1)u_{g\alpha}^p(\omega + \omega_1) + T_{wp}(s + j\omega_1)\Delta\omega_f(\omega) \\ i_{\alpha}^n(\omega - \omega_1) = T_i(s - j\omega_1)i_{\alpha r}^n(\omega - \omega_1) - Y_i(s - j\omega_1)u_{g\alpha}^n(\omega - \omega_1) + T_{wn}(s - j\omega_1)\Delta\omega_f(\omega) \end{cases} \quad (7)$$

where:

$$\begin{cases} Y_i(s) = \frac{1}{H(s)G_d(s) + Ls + r} \\ T_i(s) = H(s)G_d(s)Y_i(s) \\ T_{wp}(s) = H_p^x(s)G_d(s)Y_i(s) \\ T_{wn}(s) = H_n^x(s)G_d(s)Y_i(s) \end{cases} \quad (8)$$

3.5. Admittance Matrix

The admittance matrix of the converter, $Y_{\alpha\beta}(s)$, is defined in Equation (10):

$$\begin{bmatrix} -i_{\alpha}^p(\omega + \omega_1) \\ -i_{\alpha}^n(\omega - \omega_1) \end{bmatrix} = \begin{bmatrix} Y_{pp}(s + j\omega_1) & Y_{pn}(s + j\omega_1) \\ Y_{np}(s - j\omega_1) & Y_{nn}(s - j\omega_1) \end{bmatrix} \begin{bmatrix} u_{g\alpha}^p(\omega + \omega_1) \\ u_{g\alpha}^n(\omega - \omega_1) \end{bmatrix} = Y_{\alpha\beta}(s) \begin{bmatrix} u_{g\alpha}^p(\omega + \omega_1) \\ u_{g\alpha}^n(\omega - \omega_1) \end{bmatrix} \quad (9)$$

The elements can be derived from Figure 5, as follows:

$$\begin{cases} Y_{pp}(s + j\omega_1) = Y_i(s + j\omega_1) - (j0.5I_{dr}H_{sp}^y(s)T_i(s + j\omega_1) + T_{wp}(s + j\omega_1)H_{wp}^y(s)) \\ Y_{pn}(s + j\omega_1) = -j0.5I_{dr}T_i(s + j\omega_1)H_{sn}^y(s) - T_{wp}(s + j\omega_1)H_{wn}^y(s) \\ Y_{np}(s - j\omega_1) = j0.5I_{dr}T_i(s - j\omega_1)H_{sp}^y(s) - T_{wn}(s - j\omega_1)H_{wp}^y(s) \\ Y_{nn}(s - j\omega_1) = Y_i(s - j\omega_1) - (-j0.5I_{dr}T_i(s - j\omega_1)H_{sn}^y(s) + T_{wn}(s - j\omega_1)H_{wn}^y(s)) \end{cases} \quad (10)$$

4. Impedance-Based Stability Analysis

4.1. Admittances Analysis and Verifications

Figures 6 and 7 shows the magnitude responses of the admittances of the grid-connected converter with SRF-PLL and DSOGI-FLL, respectively. The solid lines are plotted using the theoretical models in Equation (11), while the circles are the point-by-point numerical simulation results of the admittances for comparison. The parameters used in simulations are as shown in Table 3 with the bandwidth of the SRF-PLL $f_{b\omega_PLL} = 40$ Hz (the corresponding parameters of the $G_{PLL}(s)$ is designed based on [30]), and the parameters of the DSOGI-FLL are: $k = 1.1$, $\gamma = 41$. It can be observed from Figures 6 and 7 that for different implementations of adaptive PR controllers and different PLLs, the numerical admittances match the theoretical admittances of the grid-connected converter, which verify the correctness of the proposed admittance model.

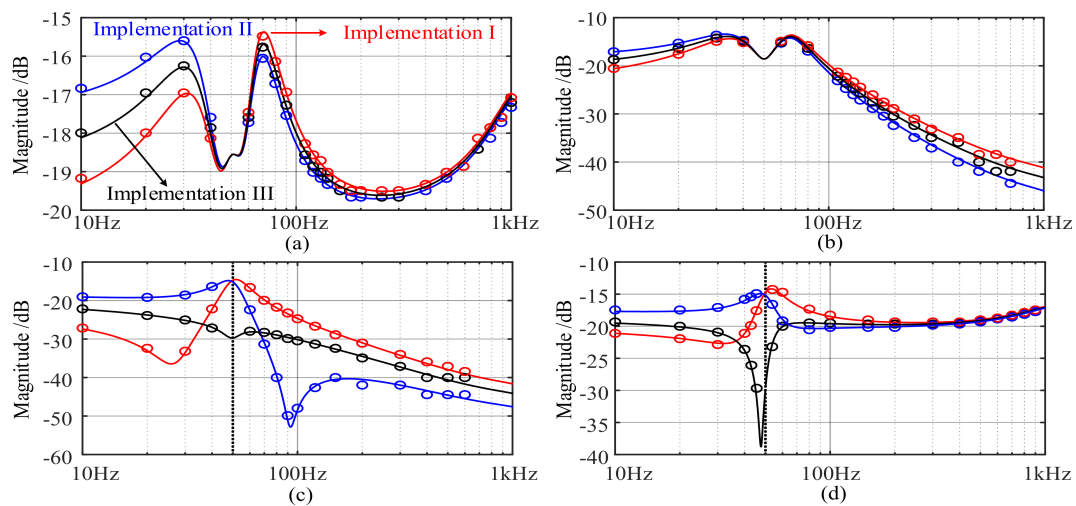


Figure 6. Magnitude responses of admittances of the converter synchronized by a SRF-PLL: (a) $Y_{pp}(s + j\omega_1)$; (b) $Y_{pm}(s + j\omega_1)$; (c) $Y_{np}(s - j\omega_1)$; (d) $Y_{m}(s - j\omega_1)$. Red line: Implementation I; Green line: implementation II; Black line: implementation III is applied. Circles: numerical simulation results.

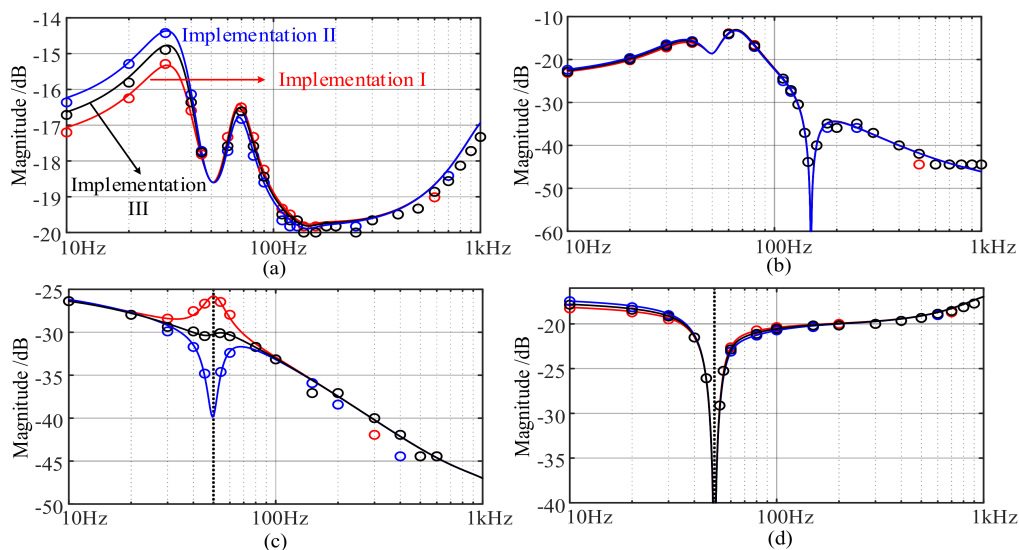


Figure 7. Magnitude responses of admittances of the converter synchronized by a DSOGI-FLL: (a) $Y_{pp}(s + j\omega_1)$; (b) $Y_{pm}(s + j\omega_1)$; (c) $Y_{np}(s - j\omega_1)$; (d) $Y_{m}(s - j\omega_1)$. Red line: Implementation I; Green line: implementation II; Black line: implementation III is applied. Circles: numerical simulation results.

Table 3. Parameters of grid-tied converter prototype.

Symbol	Description	Value
V_1	Grid phase-neutral peak voltage	$30\sqrt{2}$ V
ω_1	Grid angular frequency	$2\pi \times 50$ rad/s
f_s	Switching frequency	10 kHz
V_{dc}	Dc-link voltage	130 V
L	Inductance of the L-type filter	2 mH
r	Resistance of the filter	0.2 Ω
α_c	Current control loop bandwidth	$2\pi \times 833$ rad/s
k_i	Proportional gain of ac/dc current controller	10.47
k_r	R parameter of ac/dc current controller	1047
I_{dr}	D channel current reference of VSC	10 A

4.1.1. SRF-PLL is Used for Grid Synchronization

It can be observed from Figure 6 that the $Y_{pp}(s + j\omega_1)$ and $Y_{pn}(s + j\omega_1)$ with different implementations of the adaptive controllers are similar as shown in Figure 6a,b, while $Y_{np}(s - j\omega_1)$ and $Y_{nn}(s - j\omega_1)$ are different, especially at the frequencies around $s - j\omega_1 = j\omega_1$, i.e., $s = j2\omega_1$, as shown in Figure 6c,d. The value of $Y_{np}(s - j\omega_1)$ and $Y_{nn}(s - j\omega_1)$ at $s = j2\omega_1$ can be obtained by Equation (11), as follows:

$$\begin{cases} Y_{nn}(j\omega_1) = -Y_{np}(j\omega_1) = -\left(\frac{I_{dr}}{2} + j\frac{2\omega_1 V_n}{k_r}\right) T_{pll}(j2\omega_1) \cdots \text{I} \\ Y_{nn}(j\omega_1) = -Y_{np}(j\omega_1) = -\left(\frac{I_{dr}}{2} - j\frac{2\omega_1 V_n}{k_r}\right) T_{pll}(j2\omega_1) \cdots \text{II} \\ Y_{nn}(j\omega_1) = -Y_{np}(j\omega_1) = -\frac{I_{dr}}{2} T_{pll}(j2\omega_1) \cdots \text{III} \end{cases} \quad (12)$$

where Equations (12-I), (12-II), and (12-III) are the admittances of the converter with implementation I, II, and III of the adaptive PR controller, respectively.

Obviously, it can be obtained from Equation (12) that for a small k_r as suggested in [16], the magnitudes of $Y_{np}(s - j\omega_1)$ and $Y_{nn}(s - j\omega_1)$ of the converter using implementation III around $s = j2\omega_1$ are much smaller than those using implementation I or II.

4.1.2. DSOGI-FLL is Used for Grid Synchronization

It can be observed from Figure 7 that for different implementations of the adaptive resonant controllers, the converter admittances $Y_{pp}(s + j\omega_1)$, $Y_{pn}(s + j\omega_1)$, $Y_{nn}(s - j\omega_1)$ are similar, while $Y_{np}(s - j\omega_1)$ are different at the frequencies around $s = j2\omega_1$. Due to the effect of the zeros in $H_{wn}^{fll}(s - j\omega_1)$ and $H_{sn}^{fll}(s - j\omega_1)$ (as shown in Table 2) at $s = j2\omega_1$, it can be derived from Equation (11) that:

$$Y_{nn}(j\omega_1) = Y_{pn}(j3\omega_1) = 0 \quad (11)$$

Equation (13) is valid with different implantations of the adaptive PR controllers.

4.2. Stability Analysis

4.2.1. Stability Criterion

System stability is determined by applying general Nyquist stability criterion to the minor loop gain $L_{\alpha\beta}(s) = \mathbf{Z}_{g\alpha\beta}(s)\mathbf{Y}_{\alpha\beta}(s)$, where $\mathbf{Z}_{g\alpha\beta}(s)$ is the impedance matrix of the grid as shown in Equation (14) [28]. The system is stable if the Nyquist curves of the eigenvalues of $L_{\alpha\beta}(s)$, i.e., $\lambda_1(s)$ and $\lambda_2(s)$ as defined in Equation (15), do not encircle the critical point $(-1, j0)$, otherwise, the system is unstable [27,28]:

$$\mathbf{Z}_{g\alpha\beta}(s) = \begin{bmatrix} Z_g(s + j\omega_1) & 0 \\ 0 & Z_g(s - j\omega_1) \end{bmatrix} \quad (12)$$

$$\lambda_{1,2}(s) = 0.5 \frac{(Z_g(s + j\omega_1) Y_p(s + j\omega_1) + Z_g(s - j\omega_1) Y_n(s - j\omega_1)) \pm \sqrt{(Z_g(s + j\omega_1) Y_p(s + j\omega_1) + Z_g(s - j\omega_1) Y_n(s - j\omega_1))^2 - 4Z_g(s + j\omega_1) Z_g(s - j\omega_1) (Y_{pp}(s + j\omega_1) Y_{nn}(s - j\omega_1) - Y_{pn}(s + j\omega_1) Y_{np}(s - j\omega_1))}}{2} \quad (13)$$

4.2.2. Stability Analysis with Different PLLs

Figure 8 shows the Nyquist curves of characteristic loci of the grid system using different implementations of the adaptive current controllers and different PLLs. The parameters used in simulations are as shown in Table 3 with $f_{b\omega_PLL} = 75$ Hz, the parameters of the DSOGI-FLL are: $k = 1.1$, $\gamma = 41$, and the grid inductance $L_g = 6$ mH.

It can be observed from Figure 8a that if a DSOGI-FLL is used, the stabilities of grid-connected converters with different implementations of the adaptive PR controllers are similar. Since $Y_{pn}(j3\omega_1) = 0$ as derived in Equation (13), $Y_{pn}(s + j\omega_1) \times Y_{np}(s - j\omega_1)$ are very small around

$s = j2\omega_1$. Therefore, based on Equation (15), it can be concluded that although that $Y_{np}(s - j\omega_1)$ are different around $s = j2\omega_1$, the differences would have very limited influence on the $\lambda_1(s)$ and $\lambda_2(s)$.

It can be observed from Figure 8b that if a SRF-PLL is used, the stabilities of grid-connected converters with different implementations of the adaptive PR controllers are quite different. The converter with implementation I is unstable, while the converter with implementation II and III are stable. The stability boundaries of the grid-connected converters with different implementations of the adaptive PR controllers are obtained based on the stability criterion, as depicted in Figure 9. It can be observed from Figure 9 that the system with implementation III has the largest stable region, while the system with implementation I has the smallest stable region. When the short circuit ratio (SCR) is set to be 2.23, the maximum bandwidth to maintain system stability is around 73.3 Hz if implementation I is used, the maximum value extends to about 107.9 Hz if implementation II is applied, and the maximum value can be further raised to about 121.2 Hz if implementation III is applied. Above all, implementation III of the adaptive PR current controllers is suggested to be used for the grid connected converter using SRF-PLL for grid synchronization.

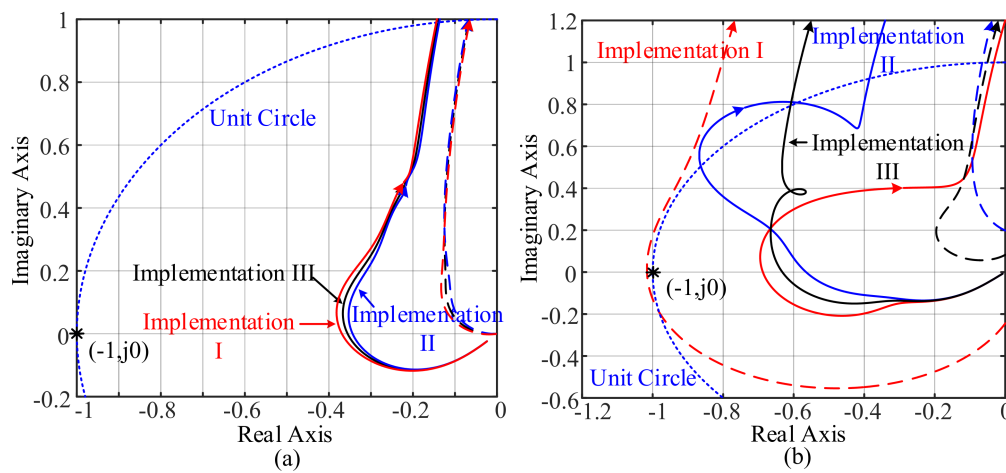


Figure 8. Nyquist curves of characteristic loci of $L_{\alpha\beta}$ with different implementations of the adaptive resonance controllers using (a) DSOGI-FLL; (b) SRF-PLL for grid synchronization. Solid line: $\lambda_1(s)$; Dash line: $\lambda_2(s)$.

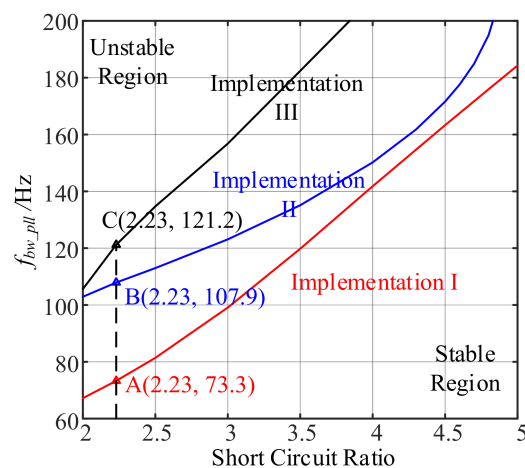


Figure 9. Stability boundaries of the grid-connected converter using different implementations of the adaptive PR controller and SRF-PLL for grid synchronization.

4.2.3. Comparison with Different PLLs

In this subsection, the effects of the SRF-PLL and DSOGI-FLL on stability of the grid connected converter with adaptive resonance current controllers using implementation III are compared.

Based on Equation (6), the dynamic properties of the output angle of a SRF-PLL with $f_{b\omega_PLL} = 40$ Hz and a DSOGI-FLL with $k = 1.1, \gamma = 41$ are similar, as shown in Figure 10a. Under this circumstance, Figure 10b shows the corresponding dynamic properties of the output frequency for different PLLs. The magnitudes of $H_{wp}^{fll}(s)$ and $H_{wn}^{fll}(s)$ are smaller than that of $H_{wp}^{pll}(s)$ and $H_{wn}^{pll}(s)$, hence it can be concluded that ω_f is more robust if the DSOGI-FLL is used than that if the SRF-PLL is used. The influence of a PLL on system stability depends on the dynamic properties of the output phase angle and the estimated frequency. Hence, the converter using DSOGI-FLL for grid synchronization should be more stable than the converter using SRF-PLL for grid synchronization.

Figure 11 shows the corresponding Nyquist curves of characteristic loci of converters with different synchronization method. The other parameters used in the simulation are as shown in Table 3. It can be concluded from Figure 11 that the converter using DSOGI-FLL for grid synchronization has a larger stability margin than that of the converter using SRF-PLL.

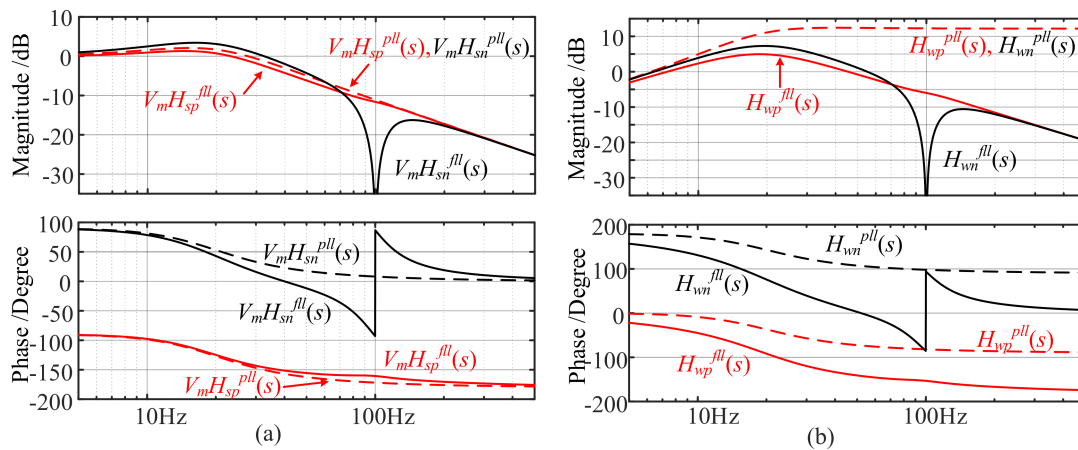


Figure 10. Frequency response of the effects of point of common coupling (PCC) voltages on θ_f in (a) and ω_f in (b). Solid line: DSOGI-FLL is used; Dash line: SRF-PLL is used.

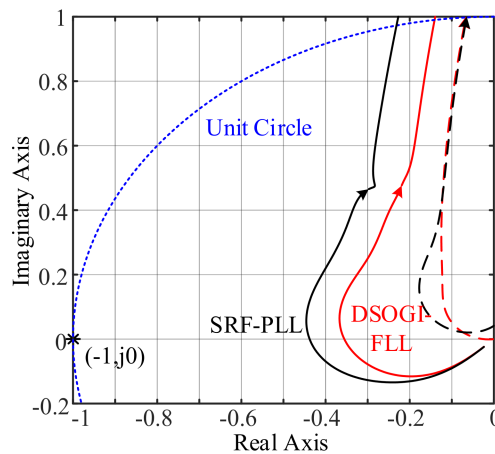


Figure 11. Nyquist curves of characteristic loci of $L_{\alpha\beta}$ using different PLLs. Black line: SRF-PLL; Red line: DSOGI-FLL is used. Solid line: $\lambda_1(s)$; Dash line: $\lambda_2(s)$.

5. Experimental Verifications

A three-phase grid-connected converter has been built and tested to verify proposed analysis. The current controllers, frame transformation and the PLL were implemented in a TMS320F28335 DSP board (Texas Instruments, Inc, Dallas, TX, USA). The grid current is sensed by a TCP0150 current probe (Tektronix, Beaverton, OR, USA) and the bandwidth of the SRF-PLL f_{bw_PLL} , the q -axis component of the grid current i_q , and the estimated frequency ω_f are sent to the D/A in the board as output signals. The output of the D/A cannot be negative, hence, it is programmed to have a 15 A offset in i_q and a 200 rad/s offset in ω_f compared to the corresponding actual values. Parameters for this experimental setup are provided in Table 3, which are consistent with the simulation parameters. The parameters of DSOGI-FLL are: $k = 1.1$, $\gamma = 41$, and the short circuit ratio is 2.23 (the corresponding L_g is 6 mH).

Figure 12 are the experimental waveforms of A-phase current of the converter using SRF-PLL for grid synchronization. In Figure 12a, the implementation I of the adaptive resonance controller is applied. At time T_0 , the bandwidth of the PLL jumps from 56 Hz to 70 Hz, and after T_0 , the grid current diverges. Once the grid currents reach the up-limited value, the system would stop running. In Figure 12b, the implementation II is applied. At time T_0 , the bandwidth of the PLL jumps from 102 Hz to 112 Hz, and after T_0 , system becomes unstable. In Figure 12c, the implementation III is applied. At time T_0 , the bandwidth of the PLL jumps from 114 Hz to 126 Hz, and after T_0 , system is no longer stable. The experimental results match the theoretical stability boundaries shown in Figure 9.

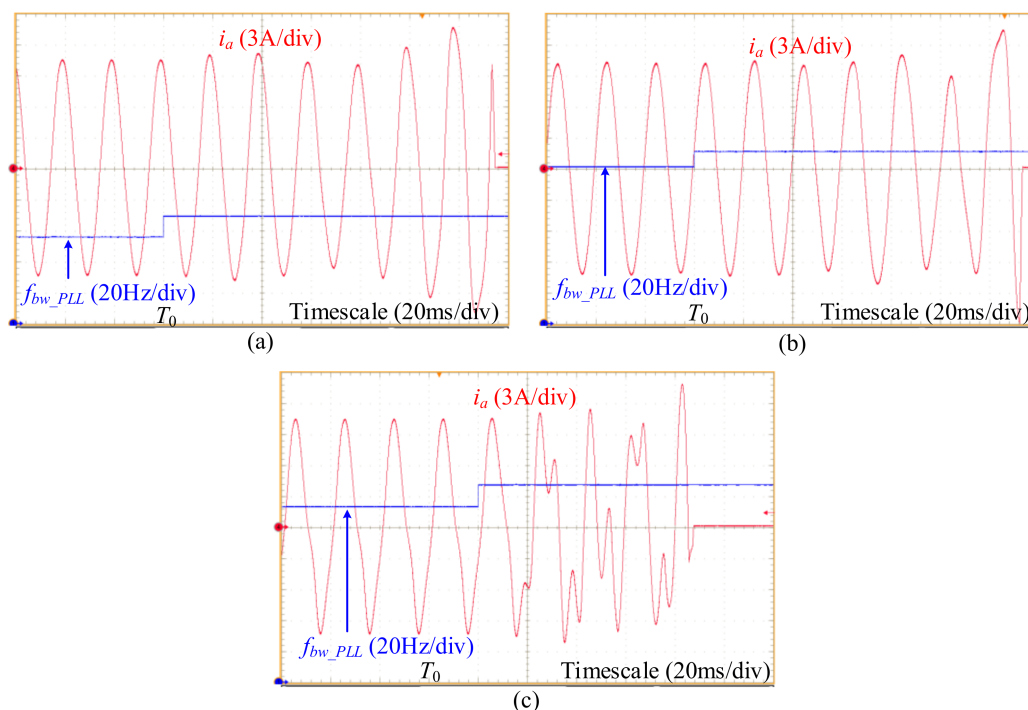


Figure 12. A-phase current waveforms for the converter using (a) implementation I; (b) implementation II; (c) implementation III of the adaptive resonance controllers with f_{bw_pll} changes at time T_0 .

Figure 13 shows the experimental waveforms of i_a , ω_f , and i_q of the grid connected converter using DSOGI-FLL for grid synchronization. In Figure 13a–c, implementation I, II and III of the adaptive PR controllers are used, respectively. The active current reference I_{dr} changes from 0 A to 10 A at time T_0 , and the dynamic responses of ω_f , and i_q are similar for all implementations of the adaptive PR controllers. The experimental results match the theoretical analysis in Section 4.2.2.

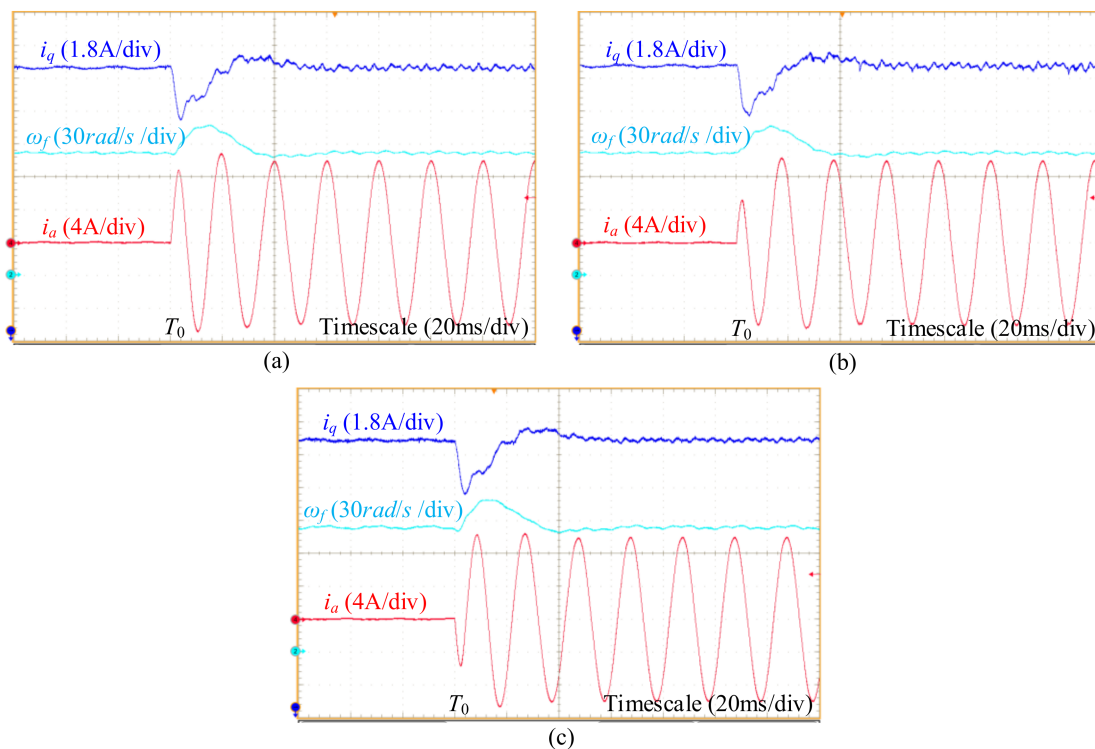


Figure 13. Dynamic responses of the converter using (a) implementation I; (b) implementation II; (c) implementation III of the adaptive resonance controllers with changing active current reference.

Figure 14 shows the experimental waveforms of A-phase current and i_q of the grid connected converter with Implementation III of the adaptive PR controllers. In Figure 14a, the SRF-PLL with 40 Hz bandwidth is applied while in Figure 14b, the DSOGI-FLL is used. The reactive current reference I_{qr} changes from -6 A to 0 A at time T_0 . It can be obtained from Figure 14 that: the percentage overshoot (PO) of i_q in Figure 14a is about 14% ($0.84/6$) while in Figure 14b the PO is less than 8% ($0.48/6$). The system with DSOGI-FLL has a stronger damping than that of the system with SRF-PLL. The experimental results verify the effectiveness of the analysis in Section 4.2.3.

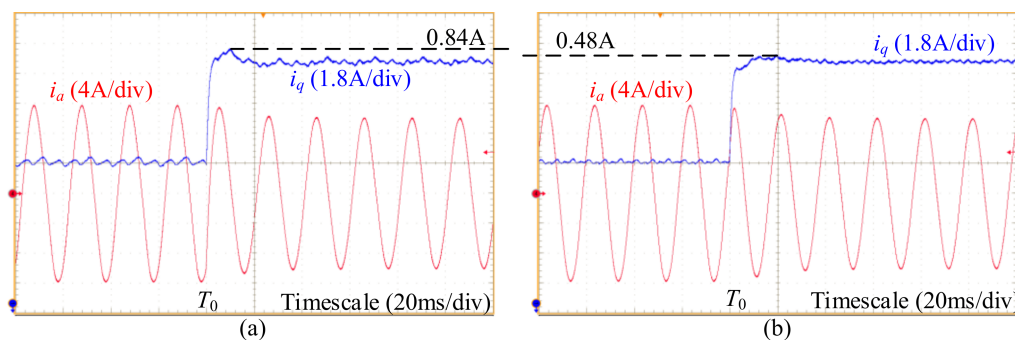


Figure 14. Dynamic response of q -axis component of grid current with its reference changes from -6 A to 0 A at time T_0 : (a) SRF-PLL; (b) DSOGI-FLL is used for grid synchronization.

6. Conclusions

The impedance model and stability of the grid-connected VSCs with adaptive resonance current controllers has been explored in this paper. Based on the proposed small-signal impedance model, some tips should be aware of when implementing adaptive resonant controllers:

- (1) If a SRF-PLL is used for grid synchronization, the system using implementation III of the resonant controller has the best stability margin, while the system using implementation I has the worst stability margin under weak grid conditions.
- (2) If a DSOGI-FLL is used for grid synchronization, the systems using implementation I, II, and III of the resonant controller have similar stability margins.
- (3) The system using a DSOGI-FLL for grid synchronization has a larger stability margin than that of the system using a SRF-PLL if the dynamic property of the output angle of the two synchronization methods are similar.

Experimental results validate the conclusions based on the theoretical analysis. In sum, the implementation of the resonance controller and the grid synchronization method should be carefully chosen for the weak-grid connected converter using adaptive resonance current controllers.

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Appendix A.

Appendix A.1. Modeling of the SRF-PLL

It can be obtained from [22] that, the output angle of the SRF-PLL is:

$$\theta_f(\omega) = -jT_{pll}(s) \left(u_{g\alpha}^p(\omega + \omega_1) - u_{g\alpha}^n(\omega - \omega_1) \right) \quad (A1)$$

where $T_{pll}(s)$ is the close loop gain of the PLL, as follows:

$$T_{pll}(s) = \frac{G_{PLL}(s)/s}{1 + V_m G_{PLL}(s)/s} \quad (A2)$$

The estimated frequency is in fact the differential of the output angle of the PLL, i.e., $\omega_f(\omega) = s\theta_f(\omega)$, hence, from Equation (A1) it can be obtained that:

$$\omega_f(\omega) = -jsT_{pll}(s) \left(u_{g\alpha}^p(\omega + \omega_1) - u_{g\alpha}^n(\omega - \omega_1) \right) \quad (A3)$$

Appendix A.2. Modeling of the DSOGI-FLL

Appendix A.2.1. Modeling of the SOGIs

According to the block diagram shown in Figure 4b, it can be obtained that:

$$\begin{cases} \varepsilon_{\alpha}^p(\omega + \omega_1) = u_{g\alpha}^p(\omega + \omega_1) - u_{\alpha}^p(\omega + \omega_1) \\ \varepsilon_{\alpha}^n(\omega - \omega_1) = u_{g\alpha}^n(\omega - \omega_1) - u_{\alpha}^n(\omega - \omega_1) \end{cases} \quad (A4)$$

The steady-state value of the u_{α} is $V_m \cos(\omega_1 t)$, by using convolution, it can be obtained that:

$$\begin{cases} \left(k\varepsilon_{\alpha}^p(\omega + \omega_1) - qu_{\alpha}^p(\omega + \omega_1) \right) \frac{\omega_1}{s+j\omega_1} - \frac{V_m}{2j} \frac{1}{s+j\omega_1} \omega_f(\omega) = u_{\alpha}^p(\omega + \omega_1) \\ \left(k\varepsilon_{\alpha}^n(\omega - \omega_1) - qu_{\alpha}^n(\omega - \omega_1) \right) \frac{\omega_1}{s-j\omega_1} + \frac{V_m}{2j} \frac{1}{s-j\omega_1} \omega_f(\omega) = u_{\alpha}^n(\omega - \omega_1) \end{cases} \quad (A5)$$

$$\begin{cases} qu_{\alpha}^p(\omega + \omega_1) = \frac{\omega_1}{s+j\omega_1} u_{\alpha}^p(\omega + \omega_1) + \frac{V_m}{j2\omega_1} \omega_f(\omega) \\ qu_{\alpha}^n(\omega - \omega_1) = \frac{\omega_1}{s-j\omega_1} u_{\alpha}^n(\omega - \omega_1) - \frac{V_m}{j2\omega_1} \omega_f(\omega) \end{cases} \quad (A6)$$

With Equations (A4)–(A6), it can be obtained that:

$$\begin{cases} u_{\alpha}^p(\omega + \omega_1) = \frac{s+j\omega_1}{D(s+j\omega_1)} \left(k\omega_1 u_{g\alpha}^p(\omega + \omega_1) + jV_m \omega_f(\omega) \right) \\ u_{\alpha}^n(\omega - \omega_1) = \frac{s-j\omega_1}{D(s-j\omega_1)} \left(k\omega_1 u_{g\alpha}^n(\omega - \omega_1) - jV_m \omega_f(\omega) \right) \end{cases} \quad (A7)$$

$$\begin{cases} qu_{\alpha}^p(\omega + \omega_1) = \frac{k\omega_1^2}{D(s+j\omega_1)} u_{g\alpha}^p(\omega + \omega_1) + \left(\frac{jV_m \omega_1}{D(s+j\omega_1)} + \frac{V_m}{j2\omega_1} \right) \omega_f(\omega) \\ qu_{\alpha}^n(\omega - \omega_1) = \frac{k\omega_1^2}{D(s-j\omega_1)} u_{g\alpha}^n(\omega - \omega_1) - \left(\frac{jV_m \omega_1}{D(s-j\omega_1)} + \frac{V_m}{j2\omega_1} \right) \omega_f(\omega) \end{cases} \quad (A8)$$

$$\begin{cases} \varepsilon_{\alpha}^p(\omega + \omega_1) = \frac{s(s+j2\omega_1)}{D(s+j\omega_1)} u_{g\alpha}^p(\omega + \omega_1) - \frac{jV_m(s+j\omega_1)}{D(s+j\omega_1)} \omega_f(\omega) \\ \varepsilon_{\alpha}^n(\omega - \omega_1) = \frac{s(s-j2\omega_1)}{D(s-j\omega_1)} u_{g\alpha}^n(\omega - \omega_1) + \frac{jV_m(s-j\omega_1)}{D(s-j\omega_1)} \omega_f(\omega) \end{cases} \quad (A9)$$

where $D(s) = s^2 + k\omega_1 s + \omega_1^2$ is used to simplify the expression.

Appendix A.2.2. Modeling of the PNSC

For a balanced positive-, negative-sequence vector, the α -, β -axis components keep the following steady-state relationship on frequency domain:

$$\begin{cases} qu_{\alpha}^p(\omega + \omega_1) = jqu_{\beta}^p(\omega + \omega_1) \\ qu_{\alpha}^n(\omega - \omega_1) = -jqu_{\beta}^n(\omega - \omega_1) \end{cases} \quad (A10)$$

Thus, the input signals of the “atan2” are:

$$\begin{cases} u_{\alpha+}^p(\omega + \omega_1) = \frac{1}{2} \frac{k\omega_1(s+j2\omega_1)}{D(s+j\omega_1)} u_{g\alpha}^p(\omega + \omega_1) + \frac{V_m}{2} \left(\frac{j(s+j2\omega_1)}{D(s+j\omega_1)} + \frac{1}{2\omega_1} \right) \omega_f(\omega) \\ u_{\alpha+}^n(\omega - \omega_1) = \frac{1}{2} \frac{k\omega_1(s-j2\omega_1)}{D(s-j\omega_1)} u_{g\alpha}^n(\omega - \omega_1) + \frac{V_m}{2} \left(\frac{-j(s-j2\omega_1)}{D(s-j\omega_1)} + \frac{1}{2\omega_1} \right) \omega_f(\omega) \end{cases} \quad (A11)$$

Appendix A.2.3. Modeling of the FLL

The steady state values of ε_{α} , ε_{β} are zeros, while qu_{α} , qu_{β} are $V_m \sin(\omega_1 t)$ and $-V_m \cos(\omega_1 t)$, respectively. Hence, by using convolution, it can easily be obtained that:

$$u_f(\omega) = -\frac{V_m}{j2} \varepsilon_{\alpha}^p(\omega + \omega_1) + \frac{V_m}{j2} \varepsilon_{\alpha}^n(\omega - \omega_1) - \frac{V_m}{2} \varepsilon_{\beta}^p(\omega + \omega_1) - \frac{V_m}{2} \varepsilon_{\beta}^n(\omega - \omega_1) \quad (A12)$$

Using the similar steady-state relationship as shown in Equation (A10), Equation (A12) can be simplified as follows:

$$u_f(\omega) = -\frac{V_m}{j} \varepsilon_{\alpha}^p(\omega + \omega_1) + \frac{V_m}{j} \varepsilon_{\alpha}^n(\omega - \omega_1) \quad (A13)$$

The steady-state value of u_f is zero, hence the output of the FLL can be written as follows:

$$-\lambda \frac{k\omega_1}{V_m} \frac{1}{s} u_f(\omega) = \omega_f(\omega) \quad (A14)$$

With Equations (A9), (A13), and (A14), the output of the FLL can be obtained that:

$$\omega_f(\omega) = \frac{-\lambda \frac{k\omega_1}{V_m} \frac{1}{j} \left(\frac{s-j2\omega_1}{D(s-j\omega_1)} u_{g\alpha}^n(\omega - \omega_1) - \frac{s+j2\omega_1}{D(s+j\omega_1)} u_{g\alpha}^p(\omega + \omega_1) \right)}{1 + \lambda \frac{k\omega_1}{s} \left(\frac{s-j\omega_1}{D(s-j\omega_1)} + \frac{s+j\omega_1}{D(s+j\omega_1)} \right)} \quad (A15)$$

Appendix A.2.4. Modeling of the Output Angle

As demonstrated in [31], the output angle of “atan2” satisfies the following relationship:

$$V_m \theta_f(\omega) = -\frac{j}{2} u_{\alpha+}^p(\omega + \omega_1) + \frac{j}{2} u_{\alpha+}^n(\omega - \omega_1) + \frac{1}{2} u_{\beta+}^p(\omega + \omega_1) + \frac{1}{2} u_{\beta+}^n(\omega - \omega_1) \quad (\text{A16})$$

It is pointed out in [27,28] that the vector at frequency $\omega - \omega_1$ is not positive-sequence, but negative-sequence, hence, Equation (A16) can be simplified to:

$$V_m \theta_f(\omega) = -j u_{\alpha+}^p(\omega + \omega_1) + j u_{\alpha+}^n(\omega - \omega_1) \quad (\text{A17})$$

With Equations (A11), (A15), and (A17), it can be obtained that:

$$\theta_f(\omega) = \frac{j \left(\frac{k\omega_1(s-j\omega_1)}{D(s-j\omega_1)} u_{\alpha+}^n(\omega - \omega_1) - \frac{k\omega_1(s+j\omega_1)}{D(s+j\omega_1)} u_{\alpha+}^p(\omega + \omega_1) \right)}{\left(1 + \frac{\lambda \left(\frac{s-j2\omega_1}{D(s-j\omega_1)} + \frac{s+j2\omega_1}{D(s+j\omega_1)} \right)}{1 + \lambda \frac{k\omega_1}{s} \left(\frac{s-j\omega_1}{D(s-j\omega_1)} + \frac{s+j\omega_1}{D(s+j\omega_1)} \right)} \right)} \times 2V_m \quad (\text{A18})$$

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